

## UNIT - I

### Electrostatics:

1. Vector Algebra
2. Coordinate System
3. Vector Calculus
4. Coulomb's Law
5. Electric field Intensity ( $E$ )
6. Electric field due to continuous charge distribution
  - Due to Point charge
  - Due to line charge
  - Due to Surface charge
  - Due to Volume charge
7. Electric flux & Electric flux Density ( $D$ )
8. Relation between  $E$  and  $D$
9. Gauss's law and its Limitation
10. Applications of Gauss's law
11. Electric Potential ( $V$ )
12. Relation between  $E$  and  $V$
13. Maxwell's Two equations for Electrostatic fields
14. Poisson's and Laplace's equations
15. Dipole and Dipole Moment
16. Electrostatic Energy and Energy Density
17. Capacitance
  - Parallel plate capacitor
  - coaxial (or) cylindrical capacitor
  - Spherical capacitor

# ELECTRO MAGNETIC FIELDS & WAVES

## Electro Magnetic Field :

It is a branch of physics in which electric & magnetic fields are studied.

- Electromagnetics (EM) may be regarded as the study of the interactions between electric charges at rest and in motion.
- EM principles find applications in various disciplines such as Antennas, Microwaves, satellite communications, fiber optics, RADARS, Radio, Television, Telephone .. etc.
- \* Vector Analysis simplifies the formulation of various laws of EM fields. Therefore the basic concepts of vector algebra, coordinate system and vector calculus should be studied before entering into the EM fields.

## 1. Vector Algebra :

It is a mathematical tool that can be used to express the quantities in different coordinate systems with magnitude and direction.

- A quantity can be either a scalar (or) a vector.

Scalar : It is a quantity that has only magnitude.

Ex Time, Temperature, mass, distance, electric potential. etc.,

A scalar is represented simply by a letter like A, B, U & V.

Vector: It is a quantity that has both magnitude & direction.

Ex Velocity, Force, displacement and electric field intensity.

A vector is represented by a letter with an arrow on top of it, such as  $\vec{A}$  and  $\vec{B}$  (or) bold letters such as A and B.

Field: Any quantity may have either scalar (or) vector field.

A field is a function that specifies a particular quantity everywhere in a region.

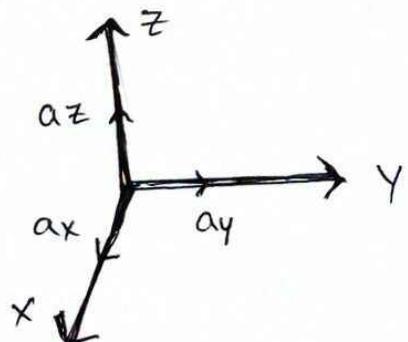
Examples of scalar fields are 1. Temperature distribution in a building  
2. Sound intensity in a theater  
3. Electric Potential in a region

Examples of vector fields are 1. Fan rotation      3. Electric forces  
2. current flow      4. Magnetic forces.

Unit Vector:

A unit vector  $\hat{a}_A$  along  $\vec{A}$  is defined as a vector whose magnitude is unity (i.e., 1) and its direction is along  $\vec{A}$ . That is

$$\hat{a}_A = \frac{\vec{A}}{|\vec{A}|} = \frac{\vec{A}}{A} .$$



\*  $\vec{A}$  is a vector and it is represented in cartesian coordinates as  $(A_x, A_y, A_z)$  (or)  $A_x a_x + A_y a_y + A_z a_z$

→ where  $A_x$ ,  $A_y$  and  $A_z$  are called components (or) magnitudes of  $A$  in the x, y & z directions respectively

→  $a_x$ ,  $a_y$  and  $a_z$  are called unit vectors in the x, y and z-directions respectively.

The magnitude of  $\vec{A}$  is given by

$$|A| = \sqrt{A_x^2 + A_y^2 + A_z^2}$$

$\therefore$  the unit vector along  $\vec{A}$  is given by

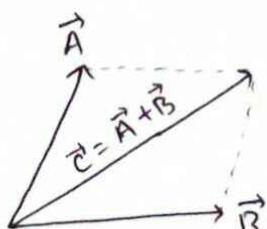
$$\vec{\alpha}_A = \frac{A_x a_x + A_y a_y + A_z a_z}{\sqrt{A_x^2 + A_y^2 + A_z^2}}$$

Here  $|\vec{\alpha}_A| = 1$ ; Hence  $\vec{A} = |A| \vec{\alpha}_A$

### Vector Addition and Subtraction :

Two vectors  $\vec{A}$  &  $\vec{B}$  can be added together to give another vector  $C$ ;

i.e.,  $\vec{C} = \vec{A} + \vec{B}$



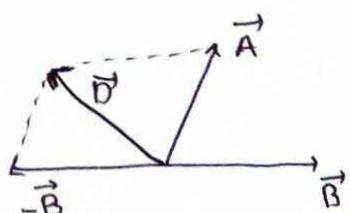
→ Vector addition is carried out component by component. Thus, if  $\vec{A} = (A_x, A_y, A_z)$  &  $\vec{B} = (B_x, B_y, B_z)$  then

$$\vec{C} = \vec{A} + \vec{B} = (A_x + B_x) a_x + (A_y + B_y) a_y + (A_z + B_z) a_z.$$

→ Vector subtraction is similarly carried out as

$$\vec{D} = \vec{A} - \vec{B} = \vec{A} + (-\vec{B})$$

$$\vec{D} = (A_x - B_x) a_x + (A_y - B_y) a_y + (A_z - B_z) a_z.$$

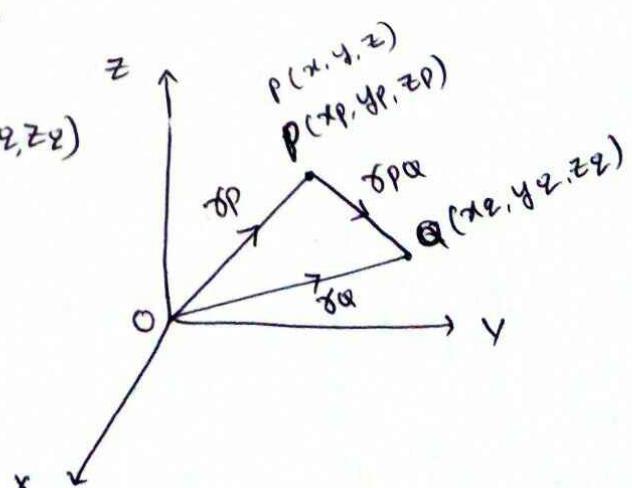


### Position and Distance Vectors :

Consider two points  $P(x_p, y_p, z_p)$  &  $Q(x_q, y_q, z_q)$

$$\vec{r}_P + \vec{r}_{PQ} - \vec{r}_Q = 0$$

$$\vec{r}_{PQ} = \vec{r}_Q - \vec{r}_P$$



→ Position vector (or) radius vector is a vector which represents the distance from the origin to the point P.

$$\therefore \vec{A} = \vec{OP} = x_a \vec{ax} + y_a \vec{ay} + z_a \vec{az}$$

∴ Here  $\vec{OP}$  is a position vector.

→ Distance vector is a vector which represents the distance from one point to another point.

$$\vec{r}_{PQ} = \vec{r}_Q - \vec{r}_P$$

∴ Here  $\vec{r}_{PQ}$  is a distance vector.

$$\vec{r}_{PQ} = (x_Q - x_P) \vec{ax} + (y_Q - y_P) \vec{ay} + (z_Q - z_P) \vec{az}$$

### Vector Multiplication :

- When two vectors A & B are multiplied, the result is either a scalar (or) a vector depending on how they are multiplied.
- There are two types of vector multiplication.
  - 1. Scalar product ( $A \cdot B$ )
  - 2. Vector product ( $A \times B$ )

1. scalar product (or) Dot product : of two vectors is defined as the product of the magnitudes of A and B and cosine of the angle between them.

$$\vec{A} \cdot \vec{B} = |A| |B| \cos \theta_{AB}$$

$$\theta_{AB} = \cos^{-1} \left( \frac{\vec{A} \cdot \vec{B}}{|A| |B|} \right)$$

→ Dot product is a scalar quantity.

If  $\vec{A} = (A_x, A_y, A_z)$  &  $\vec{B} = (B_x, B_y, B_z)$  then  $\vec{A} \cdot \vec{B} = A_x B_x + A_y B_y + A_z B_z$

Properties : i)  $A \cdot B = B \cdot A$  commutative law

ii)  $A \cdot (B + C) = A \cdot B + A \cdot C$  distributive law

$$A \cdot A = |A|^2 = A^2$$

iii)  $a_x \cdot a_x = a_y \cdot a_y = a_z \cdot a_z = 1$

$$a_x \cdot a_y = a_y \cdot a_z = a_z \cdot a_x = 0$$

## 2. Vector product (8) Cross Product

[EMF-3]

The cross product of two vectors  $\vec{A}$  &  $\vec{B}$  is defined as the magnitudes of  $A$  &  $B$  and  $\sin \theta$  of the angle between them.

$$\vec{A} \times \vec{B} = |A| |B| \sin \theta_{AB} \hat{n}$$

$\hat{n} \rightarrow$  unit vector normal (perpendicular) to the plane containing  $A$  &  $B$ .

Let  $A = (A_x, A_y, A_z)$  (8)  $A_x a_x + A_y a_y + A_z a_z$

$B = (B_x, B_y, B_z)$  (8)  $B_x a_x + B_y a_y + B_z a_z$ , then

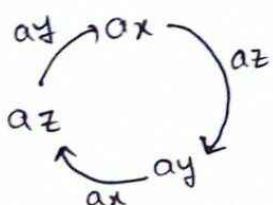
$$\vec{A} \times \vec{B} = \begin{vmatrix} a_x & a_y & a_z \\ A_x & A_y & A_z \\ B_x & B_y & B_z \end{vmatrix}$$

Properties :

i)  $A \times B = -B \times A$  anti commutative

ii)  $A \times (B+C) = A \times B + A \times C$  distributive.

$$A \times A = 0$$

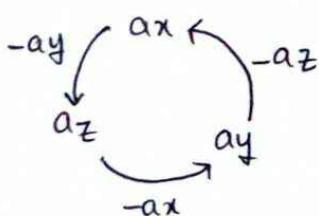


clock wise  
moving leads  
to positive

$$a_x \times a_y = a_z$$

$$a_y \times a_z = a_x$$

$$a_z \times a_x = a_y$$



anti clock wise  
moving leads  
to negative

$$a_x \times a_y = -a_z$$

$$a_y \times a_z = -a_x$$

$$a_z \times a_x = -a_y$$

$$a_x \times a_x = a_y \times a_y = a_z \times a_z = 0.$$

## Scalar Triple Product :

→ If three vectors  $\vec{A}$ ,  $\vec{B}$  &  $\vec{C}$  are given then the scalar triple product is defined as

$$A \cdot (B \times C) = B \cdot (C \times A) = C \cdot (A \times B)$$

If  $A = (A_x, A_y, A_z)$ ,  $B = (B_x, B_y, B_z)$  &  $C = (C_x, C_y, C_z)$  then

$$A \cdot (B \times C) = [ABC] = \begin{vmatrix} A_x & A_y & A_z \\ B_x & B_y & B_z \\ C_x & C_y & C_z \end{vmatrix}$$

$$B \cdot (C \times A) = [BCA]$$

$$C \cdot (A \times B) = [CAB]$$

→ Scalar triple product is a scalar quantity.

## Vector Triple product :

→ If three vectors  $\vec{A}$ ,  $\vec{B}$  &  $\vec{C}$  are given then the vector triple product is defined as

$$A \times (B \times C) = B(A \cdot C) - C(A \cdot B)$$

$$B \times (C \times A) = C(B \cdot A) - A(B \cdot C)$$

$$C \times (A \times B) = A(C \cdot B) - B(C \cdot A) \text{ and also}$$

$$(A \cdot B)C = C(A \cdot B)$$

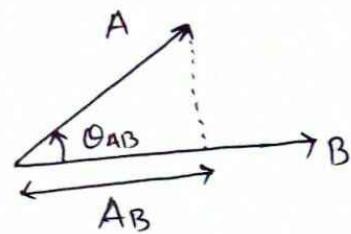
→ Vector triple product is a vector quantity.

## Components of a Vector:

EMF-4

A direct application of vector product is its use in determining the projection (or component) of a vector in a given direction. The projection can be scalar (or) vector.

Given a vector  $\vec{A}$ , we define the scalar component  $AB$  of  $\vec{A}$  along vector  $\vec{B}$  as



$$AB = \vec{A} \cos \theta_{AB} = |A| |ab| \cos \theta_{AB}$$

\* 
$$AB = \vec{A} \cdot \vec{ab}$$

The vector component  $\vec{AB}$  of  $\vec{A}$  along  $\vec{B}$  is simply the scalar component in above equation multiplied by a unit vector along  $\vec{B}$ : i.e., \* 
$$\vec{AB} = AB \vec{ab} = (\vec{A} \cdot \vec{ab}) \vec{ab}$$

### Problem

Given vectors  $A = -ax + 2ay + 2az$ ,  $B = 2ax - ay + 2az$  and  $C = 2ax + 2ay - az$ . find a)  $|A|$ ,  $|B|$  &  $|C|$  b)  $A \cdot B$ ,  $B \cdot C$  &  $C \cdot A$   
c)  $A \times B$ ,  $B \times C$  &  $C \times A$  d)  $\theta_{AB}$ ,  $\theta_{BC}$  &  $\theta_{CA}$   
e)  $A \cdot (B \times C)$ ,  $B \cdot (C \times A)$  &  $C \cdot (A \times B)$  and f)  $A \times (B \times C)$ .

Sol a)  $|A| = \sqrt{A_x^2 + A_y^2 + A_z^2} = \sqrt{1+4+4} = \sqrt{9} = 3$

$$|B| = \sqrt{4+1+4} = 3$$

$$|C| = \sqrt{4+4+1} = 3$$

b)  $A \cdot B = A_x B_x + A_y B_y + A_z B_z = (-1)(2) + 2(-1) + 2(2) = -2 - 2 + 4 = 0$

$$B \cdot C = 4 - 2 - 2 = 0$$

$$C \cdot A = -2 + 4 - 2 = 0$$

c)  $A \times B = \begin{vmatrix} ax & ay & az \\ -1 & 2 & 2 \\ 2 & -1 & 2 \end{vmatrix}$

$$= ax(4+2) - ay(-2-4) + az(+1-4)$$

$$= 6ax + 6ay - 3az$$

$$B \times C = \begin{vmatrix} ax & ay & az \\ 2 & -1 & 2 \\ 2 & 2 & 1 \end{vmatrix} = ax(1-4) - ay(-2-4) + az(4+2) \\ = -3ax + 6ay + 6az$$

$$C \times A = \begin{vmatrix} ax & ay & az \\ 2 & 2 & 1 \\ -1 & 2 & 2 \end{vmatrix} = ax(4+2) - ay(4-1) + az(4+2) \\ = 6ax - 3ay + 6az.$$

d).  $\Theta_{AB} = \cos^{-1}\left(\frac{A \cdot B}{|A||B|}\right)$   
 $= \cos^{-1}\left(\frac{0}{3 \times 3}\right)$   
 $= 90^\circ$   $\therefore \Theta_{AB} = \Theta_{BC} = \Theta_{CA} = 90^\circ.$

e)  $A \cdot (B \times C) = [ABC] = \begin{vmatrix} Ax & Ay & Az \\ Bx & By & Bz \\ Cx & Cy & Cz \end{vmatrix}$   
 $= \begin{vmatrix} -1 & 2 & 2 \\ 2 & -1 & 2 \\ 2 & 2 & 1 \end{vmatrix}$   
 $= -1(1-4) - 2(-2-4) + 2(4+2)$   
 $= 3 + 12 + 12$   
 $= 27$

$\Rightarrow A \cdot (B \times C) \quad (08)$

$$\Rightarrow A \cdot (-3ax + 6ay + 6az) = (-ax + 2ay + 2az) \cdot (-3ax + 6ay + 6az)$$
  
 $= 3 + 12 + 12$   
 $= 27.$

$$\therefore A \cdot (B \times C) = B \cdot (C \times A) = C \cdot (A \times B) = 27.$$

f).  $A \times (B \times C) = B(A \cdot C) - C(A \cdot B)$   
 $= B(0) - C(0)$   
 $= 0 - 0$   
 $= 0.$

## 2. COORDINATE SYSTEMS:

Coordinate System is defined as a system which is used to represent a point in space. There are basically three types.

1. Cartesian coordinate System
2. Cylindrical coordinate System
3. Spherical coordinate System.

- A point (or) vector can be represented in any coordinate system, which may be orthogonal.
- An orthogonal system is one in which the coordinates are mutually perpendicular.

### 1. Cartesian (or) Rectangular coordinate System:

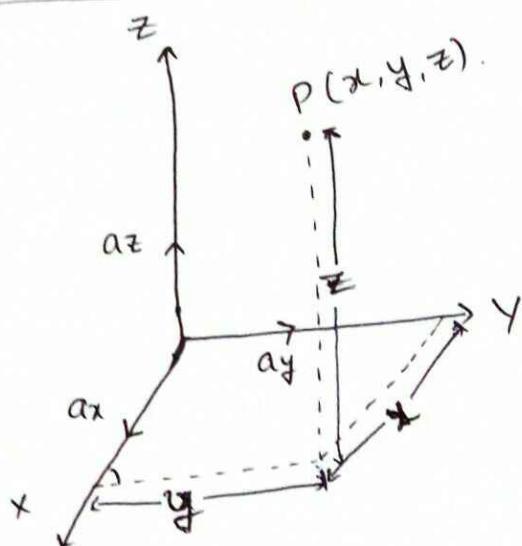
A point P can be represented as  $(x, y, z)$  shown in figure.

- The range of coordinate variables  $x, y$  and  $z$  are

$$-\infty < x < \infty$$

$$-\infty < y < \infty$$

$$-\infty < z < \infty$$



- A vector  $\vec{A}$  in cartesian coordinates can be written as  $(Ax, Ay, Az)$  (or)  $Axax + Ayay + Azaz$

where  $ax, ay$  and  $az$  are unit vectors along the  $x, y$  and  $z$  directions respectively as shown in figure. In  $P(x, y, z)$   
 ~~$x, y, z$  are distance from origin to point P along  $x, y, z$  axis respectively.~~  
 $x \rightarrow$  distance of Point P along the  $x$ -axis from the origin.

$$\begin{array}{llll} y \rightarrow & " & " & " \\ z \rightarrow & " & " & " \end{array} \quad \begin{array}{llll} \text{y-axis} & " & " \\ \text{z-axis} & " & " \end{array}$$

## 2. Cylindrical (81) Circular cylindrical coordinates : $(\rho, \phi, z)$

A point  $P$  can be represented as  $(\rho, \phi, z)$  shown in figure.

→ In figure  $\rho$  is the radius of the cylinder passing through  $P$  (or) the radial distance from the  $z$ -axis.

→  $\phi$ , called the azimuthal angle is measured from the  $x$ -axis in the  $x$ - $y$  plane.

→  $z$  is the same as in the Cartesian coordinate system.

The range of coordinate variables  $(\rho, \phi, z)$  are

$$0 \leq \rho < \infty$$

$$0 \leq \phi \leq 2\pi$$

$$-\infty < z < \infty$$

→ A vector  $\vec{A}$  in cylindrical coordinates can be written as  $(A\rho, A\phi, Az)$  (or)  $A\rho a_\rho + A\phi a_\phi + Az a_z$

where  $a_\rho$ ,  $a_\phi$  and  $a_z$  are unit vectors in the  $\rho$ ,  $\phi$  and  $z$  directions shown in figure.

→ The magnitude of  $\vec{A}$  is  $|A| = \sqrt{A\rho^2 + A\phi^2 + Az^2}$

### Properties

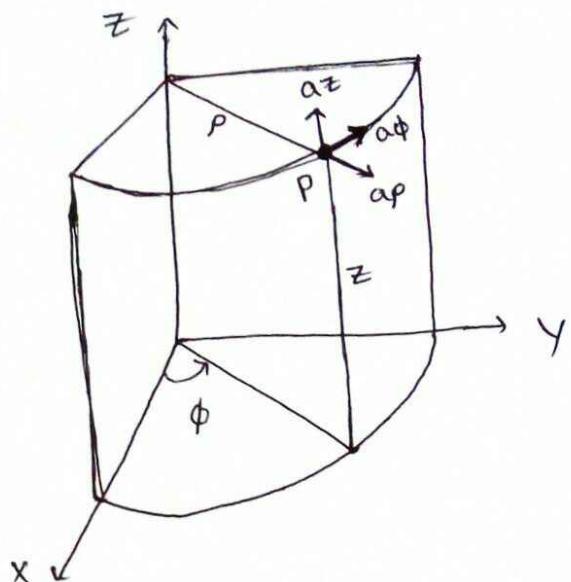
$$a_\rho \cdot a_\rho = a_\phi \cdot a_\phi = a_z \cdot a_z = 1$$

$$a_\rho \cdot a_\phi = a_\phi \cdot a_z = a_z \cdot a_\rho = 0$$

$$a_\rho \times a_\phi = a_z$$

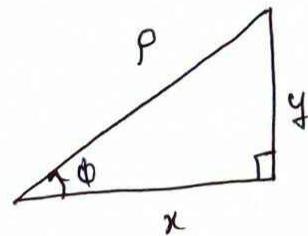
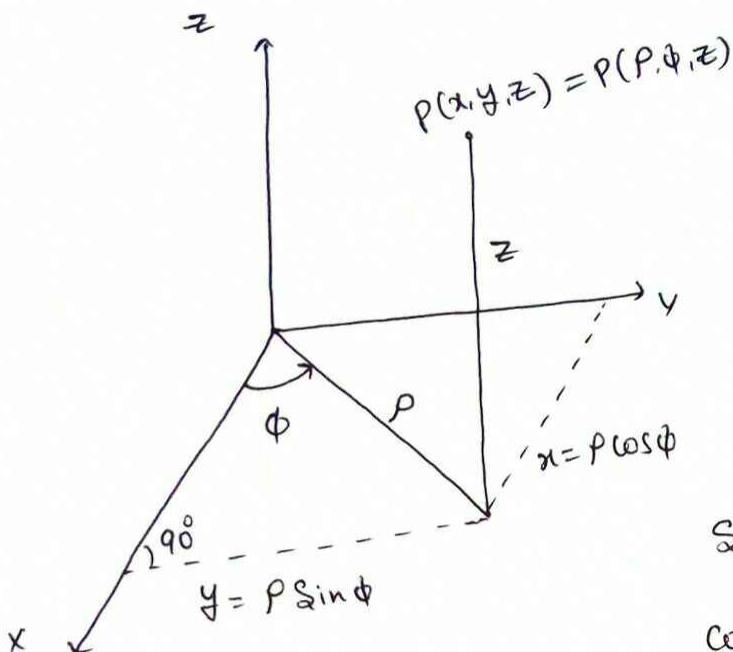
$$a_\phi \times a_z = a_\rho$$

$$a_z \times a_\rho = a_\phi$$



Relation between the variables  $(x, y, z)$  &  $(\rho, \phi, z)$ :

EMF-6



$$\sin \phi = \frac{y}{\rho} \Rightarrow y = \rho \sin \phi$$

$$\cos \phi = \frac{x}{\rho} \Rightarrow x = \rho \cos \phi$$

$$z = z$$

$$\rho^2 = x^2 + y^2 \Rightarrow \rho = \sqrt{x^2 + y^2}$$

$$\tan \phi = \frac{y}{x} \Rightarrow \phi = \tan^{-1} \frac{y}{x}$$

$$z = z$$

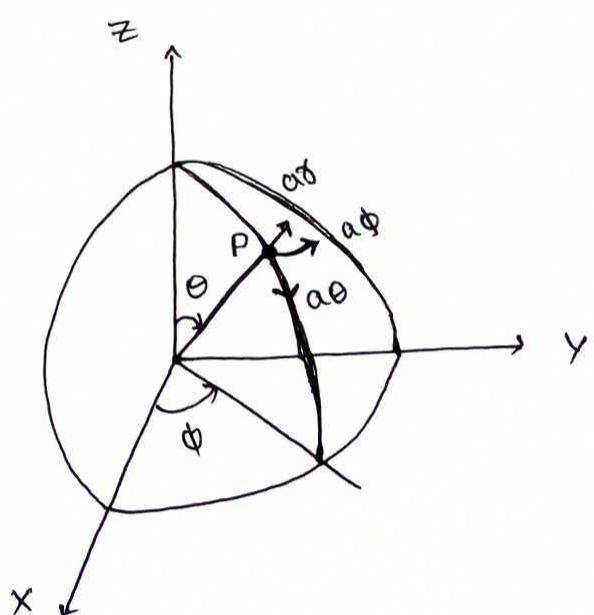
### 3. Spherical Coordinates: $(\gamma, \theta, \phi)$

A point  $P$  can be represented as  $(\gamma, \theta, \phi)$  shown in figure

where  $\gamma$  is the distance from the origin to Point P (or radius of sphere centered at the origin and passing through P).

$\theta$  called elevation angle is the angle between the  $z$ -axis and the position vector of  $P$ .

$\phi$  is same as in the cylindrical coordinate system.



→ The range of coordinate variables  $(\rho, \theta, \phi)$  are

$$0 \leq \rho < \infty$$

$$0 \leq \theta < \pi$$

$$0 \leq \phi < 2\pi$$

→ A vector  $\vec{A}$  in spherical coordinates can be written as

$$(A\rho, A\theta, A\phi) \quad (\text{or}) \quad A\rho a_\rho + A\theta a_\theta + A\phi a_\phi$$

where  $a_\rho, a_\theta$  and  $a_\phi$  are the unit vectors in  $\rho, \theta$  and  $\phi$  directions shown in figure.

→ The magnitude of  $\vec{A}$  is  $|A| = \sqrt{A\rho^2 + A\theta^2 + A\phi^2}$

Properties:

$$a_\rho \cdot a_\rho = a_\theta \cdot a_\theta = a_\phi \cdot a_\phi = 1$$

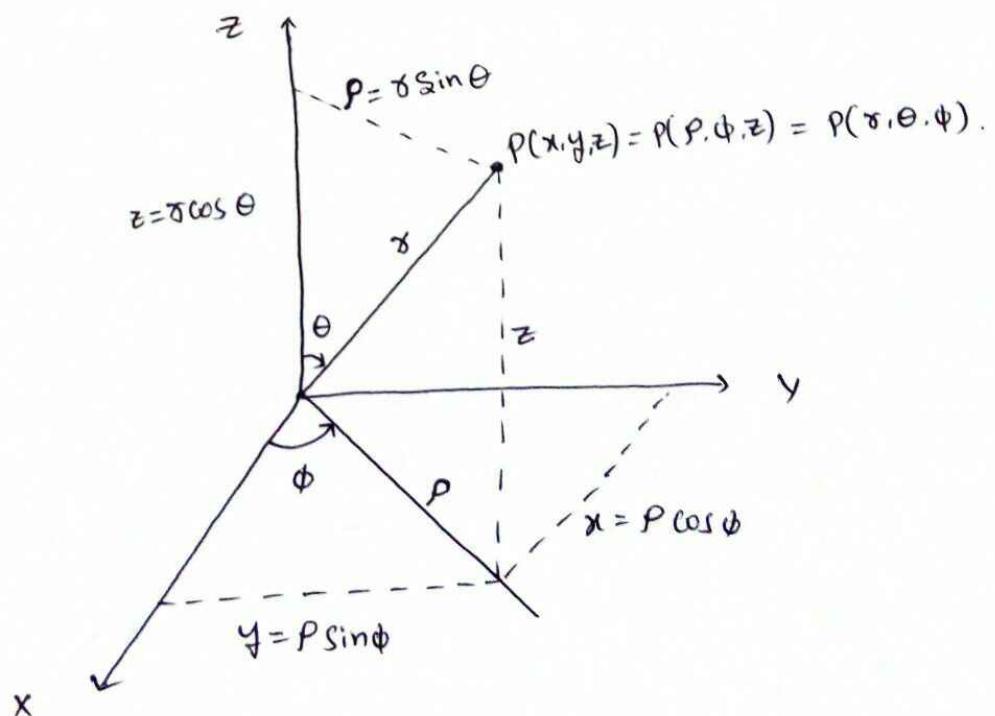
$$a_\rho \cdot a_\theta = a_\theta \cdot a_\phi = a_\phi \cdot a_\rho = 0$$

$$a_\rho \times a_\theta = a_\phi$$

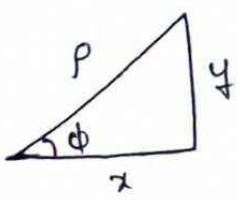
$$a_\theta \times a_\phi = a_\rho$$

$$a_\phi \times a_\rho = a_\theta$$

Relation between variables  $(x, y, z)$  &  $(\rho, \theta, \phi)$ :

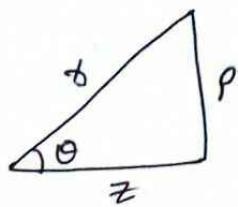


From above figure



$$\sin\phi = \frac{y}{p} \Rightarrow y = p \sin\phi$$

$$\cos\phi = \frac{x}{p} \Rightarrow x = p \cos\phi$$



$$\sin\theta = \frac{z}{p} \Rightarrow p = r \sin\theta$$

$$\cos\theta = \frac{x}{r} \Rightarrow z = r \cos\theta$$

∴ Conversion from Spherical to Cartesian System is

$$x = r \sin\theta \cos\phi$$

$$y = r \sin\theta \sin\phi$$

$$z = r \cos\theta$$

$$\Rightarrow r^2 = p^2 + z^2 = x^2 + y^2 + z^2$$

$$\therefore r = \sqrt{x^2 + y^2 + z^2}$$

$$\Rightarrow \tan\theta = \frac{p}{z} = \frac{\sqrt{x^2 + y^2}}{z}$$

$$\therefore \theta = \tan^{-1}\left(\frac{\sqrt{x^2 + y^2}}{z}\right)$$

$$\Rightarrow \phi = \tan^{-1}\left(\frac{y}{x}\right).$$

∴ Conversion from Cartesian to Spherical System is

$$r = \sqrt{x^2 + y^2 + z^2}$$

$$\theta = \tan^{-1}\left(\frac{\sqrt{x^2 + y^2}}{z}\right)$$

$$\phi = \tan^{-1}\left(\frac{y}{x}\right)$$

### Problem

Given point  $P(-2, 6, 3)$  and vector  $\vec{A} = y\alpha_x + (x+z)\alpha_y$ , express  $P$  and  $\vec{A}$  in cylindrical and spherical coordinates.

Sol

At Point  $P$ :  $x = -2$ ,  $y = 6$  &  $z = 3$ .

In cylindrical system  $r = \sqrt{x^2 + y^2} = \sqrt{4 + 36} = 6.32$

$$\phi = \tan^{-1}\left(\frac{y}{x}\right) = \tan^{-1}\left(\frac{6}{-2}\right) = 108.43^\circ$$

$$z = 3$$

In spherical System:

$$r = \sqrt{x^2 + y^2 + z^2} = \sqrt{4 + 36 + 9} = 7$$

$$\theta = \tan^{-1}\frac{\sqrt{x^2 + y^2}}{z} = \tan^{-1}\left(\frac{\sqrt{40}}{3}\right) = 64.62^\circ$$

$$\phi = \tan^{-1}\left(\frac{y}{x}\right) = \tan^{-1}\left(\frac{6}{-2}\right) = 108.43^\circ$$

$$\therefore P(-2, 6, 3) = P(6.32, 108.43^\circ, 3) = P(7, 64.62^\circ, 108.43^\circ).$$

$$A = y\alpha_x + (x+z)\alpha_y = 6\alpha_x + (-2+3)\alpha_y$$

$$\therefore A = 6\alpha_x + \alpha_y$$

### 3. VECTOR CALCULUS :

EMF-8

#### Differential Elements:

Differential elements in length, area and volume are useful in vector calculus. They are defined in the Cartesian, cylindrical and spherical coordinate systems.

#### A) Cartesian coordinates:

i) Differential length  $dl = dx \hat{ax} + dy \hat{ay} + dz \hat{az}$

ii) Differential Surface area  $dS = dy dz \hat{ax} \quad (1)$   
 $\qquad\qquad\qquad dx dz \hat{ay} \quad (2)$   
 $\qquad\qquad\qquad dx dy \hat{az}$

iii) Differential volume  $dv = dx dy dz.$

#### B) Cylindrical coordinates:

i) Differential length  $dl = d\rho \hat{ap} + \rho d\phi \hat{a\phi} + dz \hat{az}$

ii) Differential Surface

area  $dS = \rho d\phi dz \hat{ap} \quad (3) \quad d\rho dz a\phi \quad (4) \quad \rho d\rho d\phi az$

iii) Differential volume  $dv = \rho d\rho d\phi dz.$

#### C) Spherical coordinates:

i) Differential length  $dl = d\sigma \hat{a\sigma} + \tau d\theta \hat{a\theta} + \tau \sin\theta d\phi \hat{a\phi}$

ii) Differential Surface

area  $dS = \tau^2 \sin\theta d\phi d\sigma \hat{a\sigma} \quad (5)$

$\tau \sin\theta d\sigma d\phi a\theta \quad (6)$

$\tau d\sigma d\theta a\phi$

iii) Differential volume  $dv = \tau^2 \sin\theta d\sigma d\theta d\phi$

## Del ( $\nabla$ ) operator :

It is the vector differential operator and is defined as

→ In Cartesian coordinates  $\nabla = \frac{\partial}{\partial x} \mathbf{ax} + \frac{\partial}{\partial y} \mathbf{ay} + \frac{\partial}{\partial z} \mathbf{az}$

→ In cylindrical coordinates  $\nabla = \mathbf{ap} \frac{\partial}{\partial p} + \mathbf{a\phi} \frac{\partial}{\partial \phi} + \mathbf{az} \frac{\partial}{\partial z}$

→ In spherical coordinates  $\nabla = \frac{\partial}{\partial r} \mathbf{ar} + \frac{1}{r} \frac{\partial}{\partial \theta} \mathbf{a\theta} + \frac{1}{r \sin \theta} \frac{\partial}{\partial \phi} \mathbf{a\phi}$

## Gradient of a Scalar ( $\nabla V$ ) :

It is a vector quantity. It gives the maximum space-rate of change of the scalar.

Ex Temp. of soldering iron is scalar, but rate of change of Temp. is a vector.

→ In Cartesian coordinates  $\nabla V = \frac{\partial V}{\partial x} \mathbf{ax} + \frac{\partial V}{\partial y} \mathbf{ay} + \frac{\partial V}{\partial z} \mathbf{az}$

→ In cylindrical coordinates  $\nabla V = \frac{\partial V}{\partial p} \mathbf{ap} + \frac{1}{p} \frac{\partial V}{\partial \phi} \mathbf{a\phi} + \frac{\partial V}{\partial z} \mathbf{az}$

→ In spherical coordinates  $\nabla V = \frac{\partial V}{\partial r} \mathbf{ar} + \frac{1}{r} \frac{\partial V}{\partial \theta} \mathbf{a\theta} + \frac{1}{r \sin \theta} \frac{\partial V}{\partial \phi} \mathbf{a\phi}$

## Divergence of a Vector ( $\nabla \cdot A$ ) :

It is a scalar quantity. Divergence means the spreading (or) diverging of a quantity from a point.

Ex Sun rays  coming from the Sun.

→ In cartesian coordinate  $\nabla \cdot A = \frac{\partial}{\partial x} A_x + \frac{\partial}{\partial y} A_y + \frac{\partial}{\partial z} A_z$ .

→ In cylindrical coordinates  $\nabla \cdot A = \frac{\partial}{\partial p} A_p + \frac{1}{p} \frac{\partial}{\partial \phi} A_\phi + \frac{\partial}{\partial z} A_z$

→ In spherical coordinates  $\nabla \cdot A = \frac{\partial}{\partial r} A_r + \frac{1}{r} \frac{\partial}{\partial \theta} A_\theta + \frac{1}{r \sin \theta} \frac{\partial}{\partial \phi} A_\phi$

## Curl of a Vector ( $\nabla \times A$ ):

Y VNR

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It is a vector quantity. It represents the rate of rotation of a vector quantity at a point.

It is defined as circulation per unit area.

$$\rightarrow \text{In Cartesian coordinates } \nabla \times A = \begin{vmatrix} ax & ay & az \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ Ax & Ay & Az \end{vmatrix}$$

$\rightarrow$  In Cylindrical

$$\text{coordinates } \nabla \times A = \frac{1}{\rho} \begin{vmatrix} ap & p a\phi & az \\ \frac{\partial}{\partial p} & \frac{\partial}{\partial \phi} & \frac{\partial}{\partial z} \\ A\phi & PA\phi & Az \end{vmatrix}$$

$\rightarrow$  In Spherical

$$\text{coordinates } \nabla \times A = \frac{1}{\rho^2 \sin\theta} \begin{vmatrix} ar & r a\theta & r \sin\theta a\phi \\ \frac{\partial}{\partial r} & \frac{\partial}{\partial \theta} & \frac{\partial}{\partial \phi} \\ Ar & r A\theta & r \sin\theta A\phi \end{vmatrix}$$

## Laplacian of a scalar ( $\nabla^2 V$ )

$\rightarrow$  In Cartesian coordinates  $V = \nabla \cdot \nabla V = \nabla^2 V =$

$$\left[ \frac{\partial}{\partial x} ax + \frac{\partial}{\partial y} ay + \frac{\partial}{\partial z} az \right] \cdot \left[ \frac{\partial^2 V}{\partial x^2} ax + \frac{\partial^2 V}{\partial y^2} ay + \frac{\partial^2 V}{\partial z^2} az \right]$$

$$\therefore \nabla^2 V = \frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} + \frac{\partial^2 V}{\partial z^2}$$

$$\rightarrow \text{In cylindrical coordinates } \nabla^2 V = \frac{\partial^2 V}{\partial p^2} + \frac{1}{\rho^2} \frac{\partial^2 V}{\partial \phi^2} + \frac{\partial^2 V}{\partial z^2}$$

$$\rightarrow \text{In spherical coordinates } \nabla^2 V = \frac{\partial^2 V}{\partial r^2} + \frac{1}{r^2} \frac{\partial^2 V}{\partial \theta^2} + \frac{1}{r^2 \sin^2\theta} \frac{\partial^2 V}{\partial \phi^2}$$

### Stoke's Theorem :

It states that the circulation of a vector field  $\vec{A}$  around a closed path  $L$  is equal to the surface integral of the curl of  $\vec{A}$  over the open surface  $S$  bounded by  $L$

(or)

A vector relation between the line and surface integrals is called Stoke's theorem.

It is expressed as

$$\oint_L \vec{A} \cdot d\vec{l} = \int_S (\nabla \times \vec{A}) \cdot d\vec{S}$$

### Divergence Theorem :

It states that the total outward flux of a vector field  $\vec{A}$  through the closed surface  $S$  is equal to the volume integral of the divergence of  $\vec{A}$ .

(or)

The vector relation between the surface and volume integrals is called divergence theorem.

It is expressed as

$$\oint_S \vec{A} \cdot d\vec{S} = \int_{Vol} (\nabla \cdot A) dV$$

ELECTRIC FIELDCoulomb's Law:

It is an experimental law formulated in 1785 by Charles Augustin de Coulomb. This law deals with the force a point charge exerts on another point charge.

"Coulomb's law states that the force between two point charges  $Q_1$  &  $Q_2$  separated by a distance  $R$  is proportional to the product of the charges and inversely proportional to the square of the distance b/w them."

→ The direction of the force is along the line joining the charges. Coulomb also postulated that like charges repel and unlike charges attract with each other.

(or)

"Coulomb's law states that the force  $F$  between two point charges  $Q_1$  and  $Q_2$  is

- \* Along the line joining them
- \* Directly proportional to the product  $Q_1 Q_2$  of the charges.
- \* Inversely proportional to the square of the distance  $R$  b/w them."

Expressed mathematically

$$F \propto \frac{Q_1 Q_2}{R^2} = K \frac{Q_1 Q_2}{R^2}$$

Newton (N)  $\rightarrow$  (1)

where  $K \rightarrow$  proportionality constant

$$K = \frac{1}{4\pi\epsilon_0}$$

 $\epsilon_0 \rightarrow$  Permittivity of free space

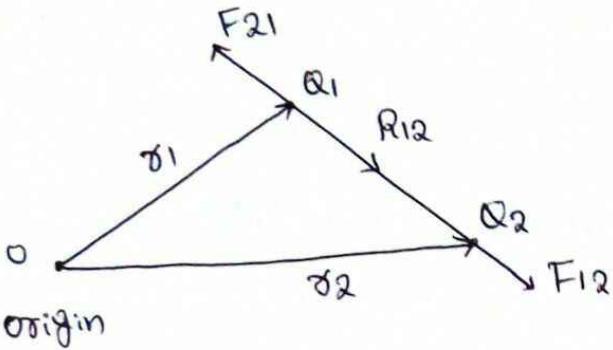
$$\therefore F = \frac{Q_1 Q_2}{4\pi\epsilon_0 R^2}$$

 $\rightarrow$  (2)

$$\epsilon_0 = 8.854 \times 10^{-12} \approx \frac{10}{36\pi} \text{ F/m}$$

$$K = \frac{1}{4\pi\epsilon_0} = 9 \times 10^9 \text{ N m/F}$$

## Coulomb vector force on point charges $Q_1$ & $Q_2$ :



→ If point charges  $Q_1$  &  $Q_2$  are located at points having position vectors  $r_1$  &  $r_2$  then the force  $F_{12}$  on  $Q_2$  due to  $Q_1$  shown in figure is given by

$$F_{12} = \frac{Q_1 Q_2}{4\pi\epsilon_0 R^2} \alpha_{R_{12}} \quad \rightarrow (3)$$

from fig  $r_1 + R_{12} - r_2 = 0 \Rightarrow R_{12} = r_2 - r_1$

and  $R = |R_{12}|$

Unit vector:  $\alpha_{R_{12}} = \frac{R_{12}}{|R_{12}|} = \frac{R_{12}}{R} \quad \rightarrow (4)$

Substituting eq (4) in (3)

$$\therefore F_{12} = \frac{Q_1 Q_2}{4\pi\epsilon_0 R^3} R_{12} = \frac{Q_1 Q_2 (r_2 - r_1)}{4\pi\epsilon_0 |r_2 - r_1|^3} \quad \rightarrow (5)$$

→ The force  $F_{21}$  on  $Q_1$  due to  $Q_2$  shown in above figure is given by  $F_{21} = |F_{12}| \alpha_{R_{21}} = |F_{12}| (-\alpha_{R_{12}})$

$\therefore F_{21} = -F_{12}$  since  $\alpha_{R_{21}} = -\alpha_{R_{12}}$ .

## Coulomb's Law:

It is an experimental law formulated in 1785 by Charles Augustin de Coulomb. This law deals with the force a point charge exerts on another point charge.

Coulomb's law states that the force  $F$  between two point charges  $Q_1$  and  $Q_2$  is

- Along the line joining them
- Directly proportional to the product  $Q_1 Q_2$  of the charges
- Inversely proportional to the square of the distance  $R$  b/w them.



$$F \propto \frac{Q_1 Q_2}{R^2}$$

$$F = K \frac{Q_1 Q_2}{R^2}$$

where  $K$  is proportionality constant

$$K = \frac{1}{4\pi\epsilon_0}$$

$\epsilon_0$  is permittivity of free space.

$$K = \frac{1}{4\pi\epsilon_0} = 9 \times 10^9 \text{ N/C}$$

$$\epsilon_0 = \frac{1}{36\pi \times 10^9} = 8.854 \times 10^{-12} \text{ C/V}$$

∴

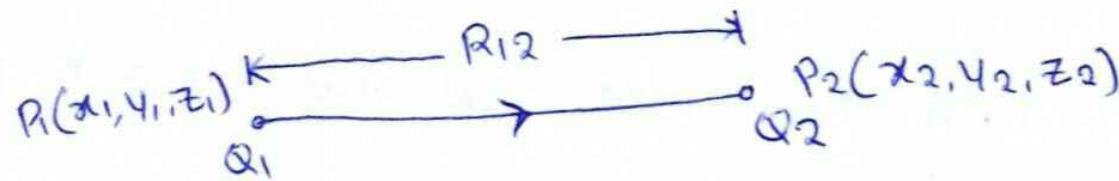
$$F = \frac{1}{4\pi\epsilon_0} \frac{Q_1 Q_2}{R^2}$$

The above force is a scalar quantity, because it has only magnitude does not have direction.

The vector force between two point charges can be computed from the formula

$$F_{12} = \frac{1}{4\pi\epsilon_0} \frac{Q_1 Q_2}{R_{12}^2} \alpha_{R_{12}}$$

where  $\alpha_{R_{12}} = \frac{R_{12}}{|R_{12}|}$  is a unit vector.



$$\mathbf{R}_{12} = (x_2 - x_1)\mathbf{ax} + (y_2 - y_1)\mathbf{ay} + (z_2 - z_1)\mathbf{az}$$

$$|\mathbf{R}_{12}| = \sqrt{(x_2 - x_1)^2 + (y_2 - y_1)^2 + (z_2 - z_1)^2}$$

→ For

point

EF-2

→ If there are more than two point charges ( $N$ ) located at position vectors  $\vec{r}_1, \vec{r}_2, \dots, \vec{r}_N$  respectively, at point with position vector  $\vec{r}$ , the resultant force  $F$  on a charge  $Q$  located at point  $\vec{r}$  is the vector sum of the forces exerted on  $Q$  by each of the charges  $Q_1, Q_2, \dots, Q_N$ . Hence

$$F = \frac{QQ_1(\vec{r}-\vec{r}_1)}{4\pi\epsilon_0|\vec{r}-\vec{r}_1|^3} + \frac{Q.Q_2(\vec{r}-\vec{r}_2)}{4\pi\epsilon_0|\vec{r}-\vec{r}_2|^3} + \dots + \frac{Q.Q_N(\vec{r}-\vec{r}_N)}{4\pi\epsilon_0|\vec{r}-\vec{r}_N|^3}$$

$$\therefore F = \frac{Q}{4\pi\epsilon_0} \sum_{k=1}^N \frac{Q_k(\vec{r}-\vec{r}_k)}{|\vec{r}-\vec{r}_k|^3}$$

### Electric Field Intensity (E) :

It is defined as the force per unit charge when placed in an electric field. It is also called as electric field strength.

E is obviously in the direction of F.

$$E = \frac{F}{Q}$$

N/C (or) Volts/meter



→ Electric field intensity at Point charge  $Q_2$  due to point charge  $Q_1$ , is given by  $E = \frac{F}{Q_2} = \frac{Q_1 Q_2}{Q_2 \cdot 4\pi\epsilon_0 R^2} \alpha R_{12} = \frac{Q_1}{4\pi\epsilon_0 R^3} R_{12}$

→ Electric field intensity at Point  $\vec{r}'$  due to point charge located at  $\vec{r}'$  is given by

$$E = \frac{Q}{4\pi\epsilon_0 R^2} \alpha R = \frac{Q(\vec{r}-\vec{r}')}{4\pi\epsilon_0 |\vec{r}-\vec{r}'|^3}$$

→ For  $N$  point charges  $Q_1, Q_2, \dots, Q_N$  located at  $\vec{r}_1, \vec{r}_2, \dots, \vec{r}_N$  the EFI at Point  $\vec{r}$  is given by  $E = \frac{Q_1(\vec{r}-\vec{r}_1)}{4\pi\epsilon_0|\vec{r}-\vec{r}_1|^3} + \frac{Q_2(\vec{r}-\vec{r}_2)}{4\pi\epsilon_0|\vec{r}-\vec{r}_2|^3} + \dots + \frac{Q_N(\vec{r}-\vec{r}_N)}{4\pi\epsilon_0|\vec{r}-\vec{r}_N|^3}$ .

$$\therefore E = \frac{1}{4\pi\epsilon_0} \sum_{k=1}^N \frac{Q_k(\vec{r}-\vec{r}_k)}{|\vec{r}-\vec{r}_k|^3}$$

Problem

Point charges  $1\text{ mC}$  &  $-2\text{ mC}$  are located at  $(3, 2, -1)$  and  $(-1, -1, 4)$ , respectively. Calculate the electric force on a  $10\text{nC}$  charge located at  $(0, 3, 1)$  and the electric field intensity at that point.

Sol

$$F = \frac{\alpha}{4\pi\epsilon_0} \sum_{k=1}^N \frac{\alpha_k \alpha R}{R^2} \hat{R}$$

Here  $k=1, 2$

$$\therefore F = \frac{\alpha}{4\pi\epsilon_0} \left[ \frac{Q_1 \alpha R_1}{R_1^2} + \frac{Q_2 \alpha R_2}{R_2^2} \right]$$

$$F = F_1 + F_2$$

$$R_1 = (0-3)\hat{x} + (3-2)\hat{y} + (1+1)\hat{z} = -3\hat{x} + \hat{y} + 2\hat{z}$$

$$|R_1| = \sqrt{9+1+4} = \sqrt{14}$$

$$R_2 = (0+1)\hat{x} + (3+1)\hat{y} + (1-4)\hat{z} = \hat{x} + 4\hat{y} - 3\hat{z}$$

$$|R_2| = \sqrt{1+16+9} = \sqrt{26}$$

$$\therefore F = \frac{\alpha}{4\pi\epsilon_0} \left[ \frac{Q_1 R_1}{|R_1|^3} + \frac{Q_2 R_2}{|R_2|^3} \right] = 10 \times 10 \times 9 \times 10^{-9} \left[ \frac{1 \times 10^{-3}(-3\hat{x} + \hat{y} + 2\hat{z})}{14\sqrt{14}} + \frac{(-2\hat{x} - 8\hat{y} + 6\hat{z})}{26\sqrt{26}} \right]$$

$$F = 9 \times 10 \times 10^{-3} \left[ \frac{-3\hat{x} + \hat{y} + 2\hat{z}}{14\sqrt{14}} + \frac{(-2\hat{x} - 8\hat{y} + 6\hat{z})}{26\sqrt{26}} \right]$$

$$F = 9 \times 10^{-2} \left[ \frac{-3\hat{x} + \hat{y} + 2\hat{z}}{14\sqrt{14}} + \frac{(-2\hat{x} - 8\hat{y} + 6\hat{z})}{26\sqrt{26}} \right]$$

$$F = -6.507\hat{x} - 3.817\hat{y} + 7.506\hat{z} \text{ mN}$$

Electric field Intensity at that point  $E = \frac{F}{\alpha}$

$$E = (-6.507\hat{x} - 3.817\hat{y} + 7.506\hat{z}) \times \frac{10^{-3}}{10 \times 10^{-9}}$$

$$E = -650.7\hat{x} - 381.7\hat{y} + 750.6\hat{z} \text{ kV/m}$$

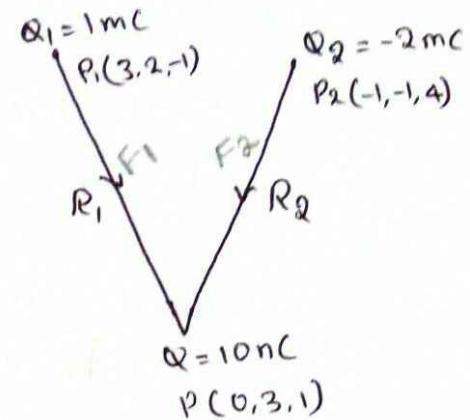
Problem

Point charges  $5\text{nC}$  &  $-2\text{nC}$  are located at  $(2, 0, 4)$  &  $(-3, 0, 5)$ , respectively.

- (a) Determine the force on a  $1\text{nC}$  point charge located at  $(1, -3, 7)$ .
- (b) Find electric field  $E$  at  $(1, -3, 7)$ .

Sol

- (a)  $-1.004\hat{x} - 1.284\hat{y} + 1.4\hat{z}$  nN & (b)  $-1.004\hat{x} - 1.284\hat{y} + 1.4\hat{z}$  V/m.



## Electric fields due to Continuous Charge Distribution :

The force and electric field charges occupying very small physical space.

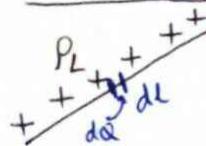
It is also possible to have continuous charge distribution along a line, on a surface and in a volume. ( $P_L$ ,  $P_S$  &  $P_V$ ).

→ The charge element  $dQ$  and the total charge  $\Omega$  due to these charge distributions are obtained from the following figures.

Point charge :



Line charge :



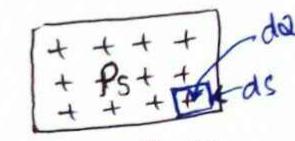
$$dQ = P_L dl$$

$$\Omega = \int_L P_L dl$$

$$E = \frac{Q}{4\pi\epsilon_0 R^2} \hat{a}_R$$

$$E = \int_L \frac{P_L dl}{4\pi\epsilon_0 R^2} \hat{a}_R$$

Surface charge :

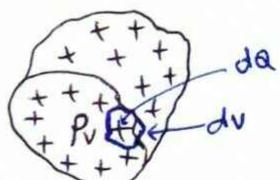


$$dQ = P_S dS$$

$$\Omega = \int_S P_S dS$$

$$E = \int_S \frac{P_S dS}{4\pi\epsilon_0 R^2} \hat{a}_R$$

Volume charge :



$$dQ = P_V dv$$

$$\Omega = \int_V P_V dv$$

$$E = \int_V \frac{P_V dv}{4\pi\epsilon_0 R^2} \hat{a}_R$$

$P_L$ ,  $P_S$  &  $P_V$  are the charge densities.

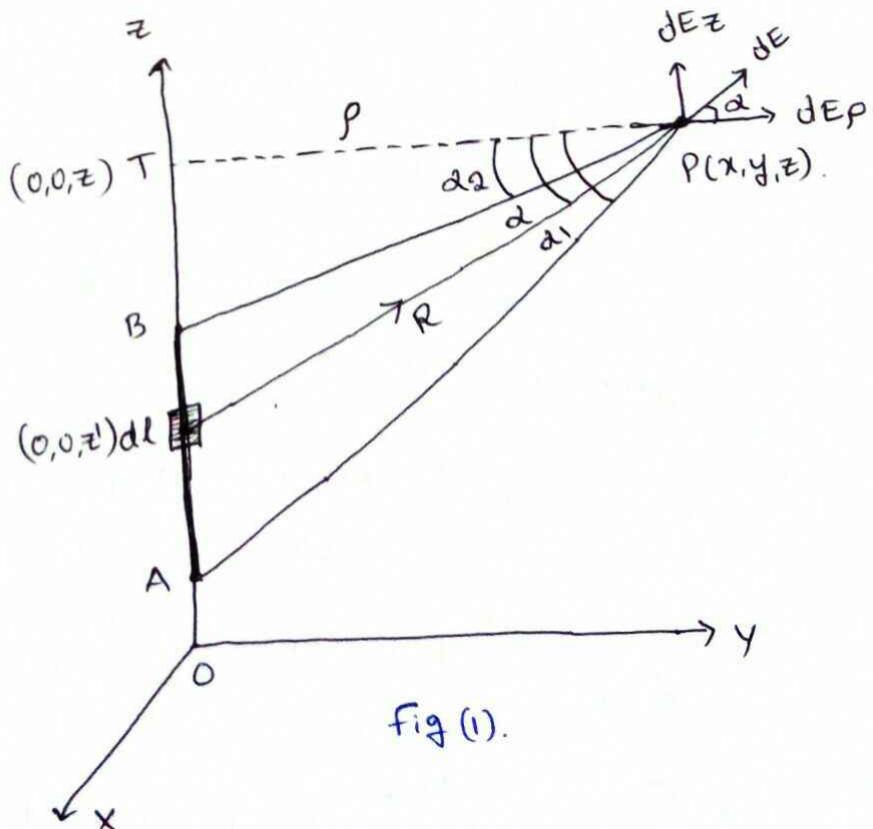
A Line charge :

→ Consider a line charge with uniform charge density  $P_L$  extended from A to B along the z-axis as shown in figure.

→ The charge element  $dQ$  associated with element  $dl = dz$  of the line is

$$dQ = P_L dl = P_L dz$$

$$\Omega = \int_{z_A}^{z_B} P_L dz$$



→ The electric field intensity  $E$  at any arbitrary point  $P(x, y, z)$

can be find by using the equation  
at  $dl = dz'$

$$E = \int_L \frac{\rho l dz'}{4\pi\epsilon_0 R^2} dR \rightarrow (1)$$

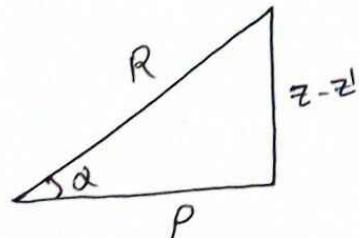
$$R = (x, y, z) - (0, 0, z')$$

from figure:

$$R = x\hat{x} + y\hat{y} + (z-z')\hat{z} \text{ (ox)}$$

$$\text{from fig(1). } R = \rho\hat{a}_\rho + (z-z')\hat{a}_z$$

$$|R| = \sqrt{\rho^2 + (z-z')^2}$$



$$\text{from (1)} E = \int_L \frac{\rho l dz'}{4\pi\epsilon_0 R^2} \cdot \frac{R}{|R|} = \int_L \frac{\rho l dz'}{4\pi\epsilon_0 R^3} R$$

$$\sin \alpha = \frac{z-z'}{R} \Rightarrow z-z' = R \sin \alpha$$

$$\cos \alpha = \frac{\rho}{R} \Rightarrow \rho = R \cos \alpha$$

$$\tan \alpha = \frac{z-z'}{\rho} \Rightarrow z-z' = \rho \tan \alpha$$



$$E = \frac{\rho l}{4\pi\epsilon_0} \int \frac{R \cos \alpha \hat{a}_\rho + R \sin \alpha \hat{a}_z}{[R^2 \cos^2 \alpha + R^2 \sin^2 \alpha]^{3/2}} dz'$$

$$E = \frac{\rho l}{4\pi\epsilon_0} \int \frac{R [\cos \alpha \hat{a}_\rho + \sin \alpha \hat{a}_z]}{[R^2 (\cos^2 \alpha + \sin^2 \alpha)]^{3/2}} dz'$$

$$\frac{R}{\rho} = \sec \alpha \Rightarrow R = \rho \sec \alpha$$

$$R^2 = \rho^2 \sec^2 \alpha$$

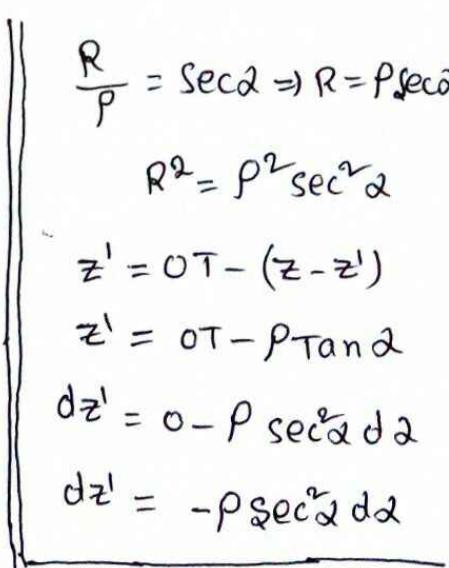
$$z' = \alpha T - (z-z')$$

$$z' = \alpha T - \rho \tan \alpha$$

$$dz' = \alpha - \rho \sec^2 \alpha d\alpha$$

$$dz' = -\rho \sec^2 \alpha d\alpha$$

$$E = \frac{\rho l}{4\pi\epsilon_0} \int \frac{(\cos \alpha \hat{a}_\rho + \sin \alpha \hat{a}_z)}{R^2} dz'$$



$$E = \frac{\rho l}{4\pi\epsilon_0} \int_{\alpha_1}^{\alpha_2} \frac{(\cos \alpha \hat{a}_\rho + \sin \alpha \hat{a}_z)}{\rho^2 \sec^2 \alpha} (-\rho \sec^2 \alpha d\alpha)$$

$$E = \frac{-\rho l}{4\pi\epsilon_0 \rho} \int_{\alpha_1}^{\alpha_2} (\cos \alpha \hat{a}_\rho + \sin \alpha \hat{a}_z) d\alpha.$$

For finite line charge,

$$E = \frac{\rho_L}{4\pi\epsilon_0 p} \left\{ [+\sin\alpha_1]_{\alpha_1}^{\alpha_2} a_p + [\cos\alpha_1]_{\alpha_1}^{\alpha_2} a_z \right\}$$

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$$E = \frac{\rho_L}{4\pi\epsilon_0 p} \left[ +(\sin\alpha_2 - \sin\alpha_1) a_p - (\cos\alpha_2 - \cos\alpha_1) a_z \right]$$

$$\boxed{E = \frac{\rho_L}{4\pi\epsilon_0 p}}$$

For an infinite line charge, Point B is at  $(0, 0, \infty)$  and A at  $(0, 0, -\infty)$  so that  $\alpha_1 = \frac{\pi}{2}$  &  $\alpha_2 = -\frac{\pi}{2}$ ; the z-vanishes

$$E = \frac{\rho_L}{4\pi\epsilon_0 p} \left[ -(\sin(-\frac{\pi}{2}) - \sin(\frac{\pi}{2})) a_p + (\cos(-\frac{\pi}{2}) - \cos(\frac{\pi}{2})) a_z \right]$$

$$E = \frac{\rho_L}{4\pi\epsilon_0 p} \left[ -(-\sin\frac{\pi}{2} - \sin\frac{\pi}{2}) a_p + (-\cos\frac{\pi}{2} - \cos\frac{\pi}{2}) a_z \right]$$

$$E = \frac{\rho_L}{4\pi\epsilon_0 p} \left[ -(-1 - 1) a_p + (-0 - 0) a_z \right]$$

$$E = \frac{\rho_L}{2 \cdot 4\pi\epsilon_0 p} [2 a_p]$$

$$\boxed{E = \frac{\rho_L}{2\pi\epsilon_0 p} a_p} \rightarrow (2)$$

Infinite line charge

$$z' = \sigma r - p \tan\alpha$$

$$(i) \alpha = \sigma r - p \tan\alpha_2$$

$$\alpha_2 = \tan^{-1}(-\infty) = -\frac{\pi}{2}$$

$$(ii) -\infty = \sigma r - p \tan\alpha_1$$

$$\alpha_1 = \tan^{-1}(\infty) = \frac{\pi}{2}$$

Eq(2) is obtained for an infinite line charge along the z-axis

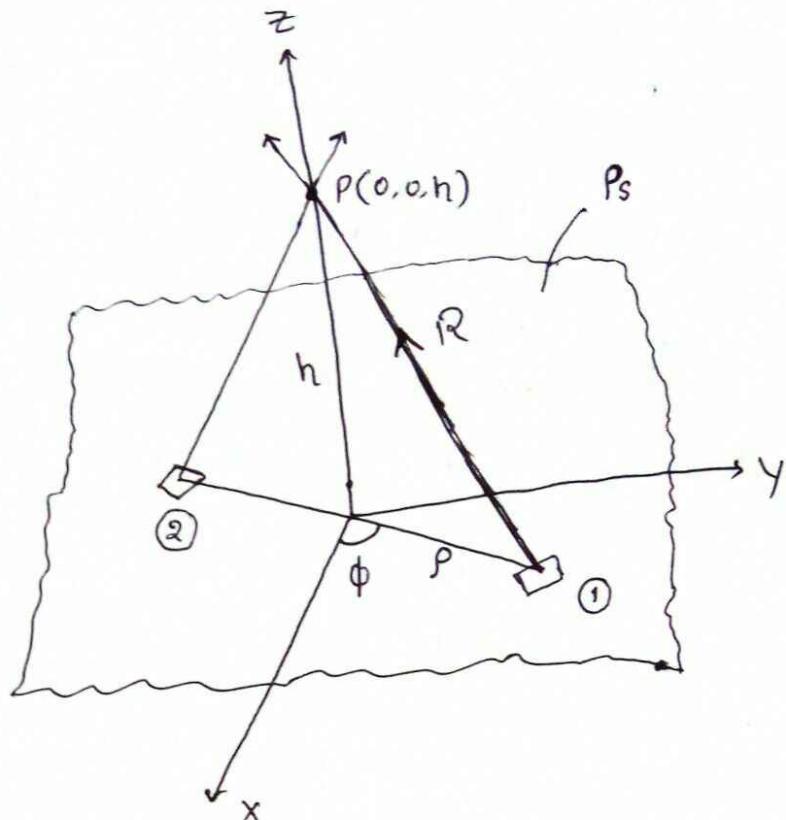
so that p & ap have their usual meaning. If the line is not along the z-axis p is the perpendicular distance from the line to the point of interest and ap is a unit vector along that distance directed from the line charge to the field point.

## A Surface Charge :

Consider an infinite sheet of charge in the x-y plane with uniform charge density  $\rho_s$ . The charge associated with an elemental area  $dS$  is

$$dQ = \rho_s dS$$

$$Q = \int_S \rho_s dS$$



→ The electric field intensity  $E$  at point  $P(0,0,h)$  by the charge  $dQ$  on the elemental surface 1. shown in figure.

$$E = \int \frac{dQ}{4\pi\epsilon_0 R^2} \cdot \hat{R} = \int \frac{\rho_s dS}{4\pi\epsilon_0 R^2} \cdot \hat{R} \quad (8) \quad dE = \frac{dQ}{4\pi\epsilon_0 R^2} \cdot \hat{R}$$

$$\text{from fig. } \rho + R - h = 0 \Rightarrow R = -\rho + h$$

$$\therefore R = -\rho \hat{a}_\rho + h \hat{a}_z$$

$$|R| = \sqrt{\rho^2 + h^2} \quad \& \quad dR = \frac{R}{|R|}$$

$$dS = \rho d\rho d\phi \quad [\because \text{Surface area along } z\text{-axis in cylindrical coordinate system}]$$

Substitute these values in the above equation.

$$\therefore E = \int \frac{\rho_s dS \cdot R}{4\pi\epsilon_0 |R|^3} = \int \frac{\rho_s \rho d\rho d\phi \cdot (-\rho \hat{a}_\rho + h \hat{a}_z)}{4\pi\epsilon_0 (\sqrt{\rho^2 + h^2})^3}$$

$$E = \frac{\rho_s}{4\pi\epsilon_0} \int \frac{(-\rho \hat{a}_\rho + h \hat{a}_z) \rho d\rho d\phi}{(\rho^2 + h^2)^{3/2}}$$

→ Due to the symmetry of the charge distribution, for every element 1, there is a corresponding element 2, whose contribution along  $A_P$  cancels that of element 1. is shown in figure.

Thus the contribution to  $E_P$  add up to zero so that  $E$  has only  $z$ -component.

$$\therefore E = \frac{\rho_s}{4\pi\epsilon_0} \int_{\phi=0}^{2\pi} \int_{\rho=0}^{\infty} \frac{h\rho d\rho d\phi}{(\rho^2 + h^2)^{3/2}} az$$

$$E = \frac{\rho_s h}{4\pi\epsilon_0} \int_0^{2\pi} \int_0^{\infty} \frac{\rho d\rho}{(\rho^2 + h^2)^{3/2}} az$$

$$E = \frac{\rho_s h}{2\epsilon_0} \int_0^{\pi/2} \frac{h \tan\theta h \sec^2\theta d\theta}{(h^2 \tan^2\theta + h^2)^{3/2}} az$$

$$E = \frac{\rho_s h}{2\epsilon_0} \int_0^{\pi/2} \frac{h^2 \tan\theta \sec^2\theta d\theta}{[h^2 (\tan^2\theta + 1)]^{3/2}} az$$

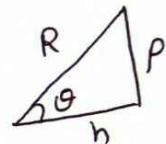
$$E = \frac{\rho_s h}{2\epsilon_0} \int_0^{\pi/2} \frac{h^2 \tan\theta \sec^2\theta d\theta}{h^3 \sec^3\theta} az$$

$$E = \frac{\rho_s}{2\epsilon_0} \int_0^{\pi/2} \frac{\tan\theta d\theta}{\sec\theta} az = \int_0^\pi \frac{\rho_s}{2\epsilon_0} \frac{\sin\theta}{\cos\theta} \cdot \cos\theta d\theta az$$

$$E = \frac{\rho_s}{2\epsilon_0} \int_0^{\pi/2} \sin\theta d\theta az = \frac{\rho_s}{2\epsilon_0} [-\cos\theta]_0^{\pi/2} = \frac{\rho_s}{2\epsilon_0} (-(-1))az = \frac{\rho_s}{2\epsilon_0} az.$$

$$\therefore E = \frac{\rho_s}{2\epsilon_0} az$$

from fig



$$\tan\theta = P/h$$

$$P = h \tan\theta$$

$$dP = h \sec^2\theta d\theta$$

$$\text{when } \rho=0; \tan\theta=0 \\ \theta=\tan^{-1}(0)=0$$

$$\text{when } \rho=\infty; \tan\theta=\infty \\ \theta=\tan^{-1}(\infty)=\frac{\pi}{2}$$

$E$  has only  $z$ -component if the charge is in the  $x-y$  plane.

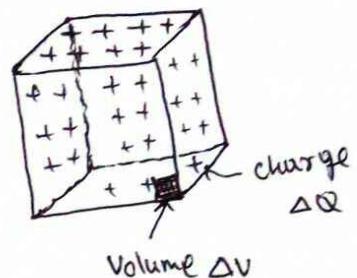
→ In general, for an infinite sheet of charge

$$E = \frac{\rho_s}{2\epsilon_0} an$$

## A Volume charge :

A charge uniformly distributed throughout a volume is called a volume charge. The volume charge density  $\rho_V$  is defined as the charge per unit volume.

$$\rho_V = \frac{Q}{V}$$



When charge  $Q$  is distributed throughout a specified volume  $V$ , then each element contributes electric field with respect to some external point.

Consider a small amount of charge  $\Delta Q$  over a small-volume  $\Delta V$ , the ratio  $\frac{\Delta Q}{\Delta V}$  is called volume charge density. It is denoted with  $\rho_V$  & measured in  $C/m^3$ .

$$\rho_V = \frac{\Delta Q}{\Delta V} = \frac{dQ}{dV} \Rightarrow dQ = \rho_V dV$$

$$\Rightarrow Q = \int \rho_V dV$$

→ Electric field induced due to a point charge is

$$E = \frac{1}{4\pi\epsilon_0} \frac{Q}{R^2} AR$$

→ Electric field induced due to a small charge is

$$dE = \frac{1}{4\pi\epsilon_0} \frac{dQ}{R^2} AR = \frac{1}{4\pi\epsilon_0} \rho_V dV \cdot AR$$

$$E = \frac{1}{4\pi\epsilon_0} \int \rho_V dV \cdot AR$$

\* A small charge  $Q$  occupying very small area (or) volume is called a point charge.

\* A charge uniformly distributed along a line is called line charge.

\* A charge uniformly distributed over a surface (or) sheet is called surface charge.

\* A charge uniformly distributed throughout a volume is called volume charge.

Problem

A circular ring of radius  $a$  carries a uniform charge  $\rho_L$  C/m and is placed on the  $xy$ -plane with axis same as the  $z$ -axis.

(a) Show that

$$E(0,0,h) = \frac{\rho_L ah}{2\epsilon_0 [h^2 + a^2]^{3/2}} \hat{az}$$

(b) What values of  $h$  gives the maximum value of  $E$ ?

(c) If the total charge on the ring is  $Q$ , find  $E$  as  $a \rightarrow 0$ .

Sol

(a) For a line charge the electric field intensity is given by

$$E = \int \frac{\rho_L dl}{4\pi\epsilon_0 R^2} \hat{ar}$$

$$\text{from fig: } R + a - h = 0$$

$$R = -a + h$$

$$\therefore R = a(-ap) + ha\hat{z}$$

$$|R| = \sqrt{a^2 + h^2}$$

$$ar = \frac{R}{|R|} \quad \& \quad dl = ad\phi.$$

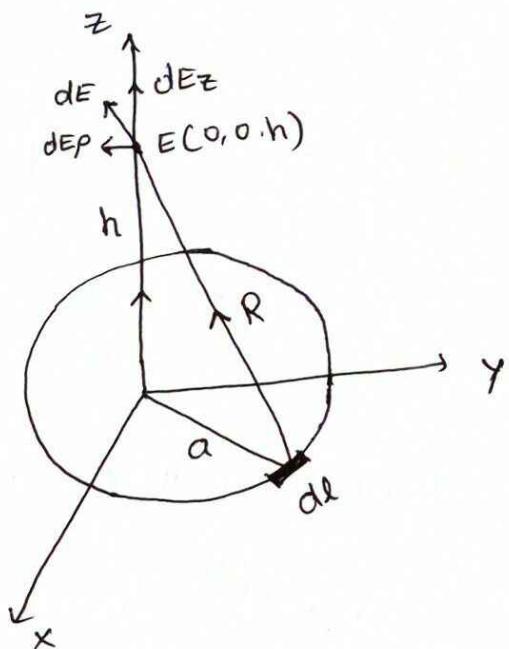
Substitute all these values in the above equation.

$$E = \int \frac{\rho_L}{4\pi\epsilon_0} \frac{(-aap + ha\hat{z})}{(\sqrt{a^2 + h^2})^3} a d\phi$$

$$E = \frac{\rho_L}{4\pi\epsilon_0} \int_{\phi=0}^{2\pi} \frac{(-aap + ha\hat{z})}{(a^2 + h^2)^{3/2}} a d\phi$$

→ By symmetry, the contribution along  $ap$  added up to zero.

This is for every element  $dl$  there is a corresponding element diametrically opposite to it. So that the two contributions cancel each other. Thus  $E$  has only  $z$ -component.



that is,  $E = \frac{\rho_L a h a_z}{4\pi\epsilon_0 (a^2 + h^2)^{3/2}} \int_0^{2\pi} d\theta = \frac{\rho_L a h a_z}{4\pi\epsilon_0 (a^2 + h^2)^{3/2}} \times 2\pi$

$$\therefore E(0,0,h) = \frac{\rho_L a h}{2\epsilon_0 (h^2 + a^2)^{3/2}}$$

(b)

$$\begin{aligned} \frac{d|E|}{dh} &= \frac{\rho_L a}{2\epsilon_0} a_z \frac{d}{dh} \frac{h}{(h^2 + a^2)^{3/2}} \\ &= \frac{\rho_L a}{2\epsilon_0} a_z \left[ \frac{(h^2 + a^2)^{3/2} \cdot (1) - h \cdot \frac{3}{2}(h^2 + a^2)^{1/2} \cdot 2h}{((h^2 + a^2)^{3/2})^2} \right] \\ &= \frac{\rho_L a}{2\epsilon_0} a_z \left[ \frac{(h^2 + a^2)^{3/2} - 3h^2(h^2 + a^2)^{1/2}}{(h^2 + a^2)^3} \right] \end{aligned}$$

For maximum value of  $E$   $\frac{d|E|}{dh} = 0$  i.e.

$$(h^2 + a^2)^{3/2} = 3h^2(h^2 + a^2)^{1/2} \Rightarrow h^2 + a^2 = 3h^2 \Rightarrow a^2 = 2h^2$$

$$\therefore h^2 = \frac{a^2}{2} \Rightarrow h = \pm \frac{a}{\sqrt{2}}$$

(c) since the charge is uniformly distributed, the line charge density is  $\rho_L = \frac{Q}{l} = \frac{\alpha}{2\pi a}$

$$E = \frac{\rho_L a h}{2\epsilon_0 (h^2 + a^2)^{3/2}} a_z = \frac{\alpha a h}{2\pi a \cdot 2\epsilon_0 (h^2 + a^2)^{3/2}} a_z$$

$$E = \frac{\alpha h}{4\pi\epsilon_0 (h^2 + a^2)^{3/2}} a_z$$

$$\text{As } a \rightarrow 0 : E = \frac{\alpha h}{4\pi\epsilon_0 h^{3/2}} a_z = \frac{\alpha}{4\pi\epsilon_0 h^2} a_z$$

In general

$$E = \frac{\alpha}{4\pi\epsilon_0 R^2} a_z$$

## Electric Flux & Flux Density:

Consider two point charges placed in a space. If one charge is brought near another charge, a force is exerted b/w them. An experiment conducted by Michael Faraday in 1837 showed that if a positive charge is placed near a negative charge, due to the force, an electric field is developed between the charges. The lines of force originated radially from positive charge and terminated at negative charge as shown in below figure.

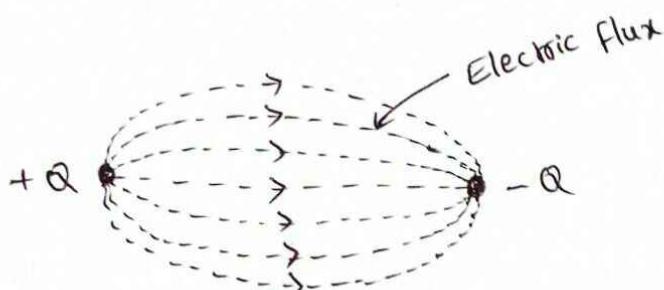


fig: Electric field  
b/w 2 Point charges.

### Electric flux ( $\Psi$ ):

The total number of lines of force in an electric field is called electric flux. The lines of force are called flux lines (or) Stream Lines.

The electric charge produces the electric flux and as the charge increases the electric flux also increases. Therefore the electric flux is numerically equal to the electric charge. If an electric charge produces  $Q$  coulombs, then the electric flux is given

by

$$\boxed{\Psi = Q \text{ coulombs}}$$

#### Note:

1. The flux lines are always parallel to each other
2. The flux lines do not depend on the medium in which the charges are placed.
3. The flux lines radiate in all directions from/into the charges.

## Electric Flux Density ( $\vec{D}$ ):

Consider the electric flux  $\Psi$  generated from a coulombs of electric charges as shown in figure.

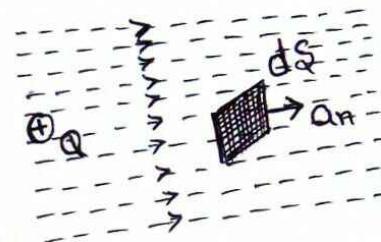


fig: Electric flux lines.

→ let the differential Surface area  $dS$  at point perpendicular to the lines of force. The electric flux crossing through this Surface-area is  $d\Psi$ .

\* Def: The net electric flux passing through the unit Surface-area is called electric flux density  $\vec{D}$ .

→ The direction of electric flux density is normal to the Surface-area. It is expressed as

$$\vec{D} = \frac{d\Psi}{dS} \vec{an}$$

c/m<sup>2</sup>

## Electric flux density (EFD) due to Point charge Q :

→ Consider an imaginary sphere of radius  $r$  and a point charge  $+Q$  is placed at the centre of the sphere as show in figure.

→ The flux lines originates from the point charge distributed radially over the Surface of the sphere.

The flux density at the differential Surface  $dS$  is

$$D = \frac{d\Psi}{dS} \vec{an}$$

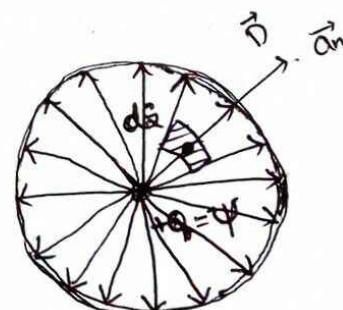


fig: flux from point charge

→ If the total Surface area is considered through which the flux  $\Psi$  is passing, the flux density is  $D = \frac{\text{Total Flux}}{\text{Surface area of sphere}}$

$$D = \frac{\Psi}{4\pi r^2} \vec{an} = \frac{Q}{4\pi r^2} \vec{an}$$

Since  $A_n = A_s$ , radial direction normal to the surface area.

then

$$\vec{D} = \frac{Q}{4\pi\epsilon_0 r^2} A_s$$

The electric flux density is also called displacement flux density.

Relation b/w Electric field Intensity & Electric flux Density :

Electric field intensity  $\vec{E} = \frac{Q}{4\pi\epsilon_0 r^2} A_s$

$$E\epsilon_0 = \frac{Q}{4\pi r^2} A_s \rightarrow (1)$$

Electric flux density  $\vec{D} = \frac{Q}{4\pi r^2} A_s \rightarrow (2)$

Equation eq (1) & (2) then  $\vec{D} = \vec{E}\epsilon_0$

$\therefore \vec{D}$  &  $\vec{E}$  are related through the permittivity of the medium.

Electric field Density Due to Line charge :

→ for infinite line charge density  $\rho_L$  C/m

we know that  $E = \frac{\rho_L}{2\pi\epsilon_0 r} A_p$

then  $D = \epsilon_0 E = \epsilon_0 \frac{\rho_L}{2\pi\epsilon_0 r} A_p = \frac{\rho_L}{2\pi r} A_p$

$\therefore$  The flux density is

$$\vec{D} = \frac{\rho_L}{2\pi r} A_p$$

## Electric flux density due to Surface (Sheet) charge:

→ for an infinite sheet charge density  $\rho_s$   $C/m^2$

we know that  $E = \frac{\rho_s}{2\epsilon_0} a_n$

then  $D = \epsilon_0 E = \epsilon_0 \frac{\rho_s}{2\epsilon_0} a_n = \frac{\rho_s}{2} a_n$ .

∴ the flux density is 
$$\vec{D} = \frac{\rho_s}{2} a_n$$

## Electric flux density due to volume charge:

→ for a volume charge density  $\rho_v$   $C/m^3$

we know that  $E = \frac{1}{4\pi\epsilon_0} \int_V \rho_v dv a_\theta$

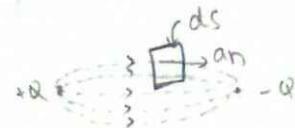
then  $D = \epsilon_0 E = \epsilon_0 \cdot \frac{1}{4\pi\epsilon_0} \int_V \rho_v dv a_\theta$

the flux density 
$$D = \frac{1}{4\pi\theta^2} \int_V \rho_v dv a_\theta$$

Gauss's Law :

Gauss's law states that the total electric flux  $\Psi$  through any closed surface is equal to the total charge enclosed by that surface.

$$\nabla \cdot D = P_V$$



$$dS = dS \cos \theta$$

$$D = \frac{d\Psi}{dS} \Rightarrow d\Psi = D dS \cos \theta = D \cdot dS$$

$D$  is normal to the surface

$dS$  direction is normal to the surface

$\therefore$  angle b/w  $D$  &  $dS$  is zero;  $\theta = 0$

i.e.  $D$  &  $dS$  are parallel

Hence  $d\Psi = D dS = D \cdot dS$

$$d\Psi = D \cdot dS \Rightarrow \Psi = \oint_S D \cdot dS$$

→ Applying divergence theorem to the above equation

$$\Psi = \int_S D \cdot dS = \int_V \nabla \cdot D \, dV$$

→ Total charge enclosed by the volume is

$$Q = \int_V P_V \, dV$$

∴ from the statement

$$\Psi = Q = \int_S D \cdot dS = \int_V \nabla \cdot D \, dV = \int_V P_V \, dV$$

It is the integral form of Gauss's law

→ The point form (a) Differential form of Gauss law is

$$\nabla \cdot D = P_V$$

→ The above equation states that the volume charge density is same as the divergence of the electric flux density.

\* Gauss's law is used for only symmetrical charge distributions.

It is the limitation of Gauss's law.

## Applications of Gauss's law :

- Gauss's law is used to compute electric field intensity ( $E$ ) and electric flux density ( $D$ ) due to symmetric charge distributions.
- \* The Gaussian surface is always a closed surface. It may be of any shape.
  - \* The surface is chosen such that at each point on the Gaussian surface, flux density  $D$  is either normal ( $\theta = 90^\circ$ ) tangential to the surface.
- ⇒ when  $D$  is normal to the Gaussian Surface 
$$D \cdot dS = D dS$$
- because  $D$  is constant on the surface
- ⇒ when  $D$  is tangential 
$$D \cdot dS = 0.$$

### 1. Point charge :

→ Suppose a point charge  $Q$  is located at the origin. To determine  $D$  at a point  $P$ , it is easy to see that choosing a spherical surface containing  $P$  will satisfy symmetry conditions.

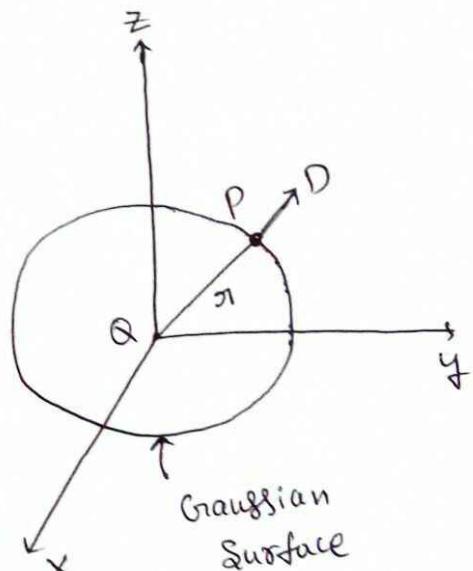
→ Thus a spherical surface centered at the origin is the Gaussian surface is shown in figure.

→ Since  $D$  is everywhere normal to the Gaussian surface, i.e.,

$$D = D \hat{\theta} \alpha_{\theta}$$

applying Gauss's law

$$\Psi = Q_{\text{enclosed}} = \oint_S D \cdot dS = D \oint_S dS$$



$dS \rightarrow$  Surface [charge] area along  $\theta$ -direction in spherical coordinate system.

$$dS = r^2 \sin\theta d\theta d\phi d\sigma \quad \& \quad D = D \hat{\theta} \alpha_{\theta} ; \oint_S D \cdot dS = \int_S D \hat{\theta} \cdot r^2 \sin\theta d\theta d\phi (\alpha_{\theta} \cdot \alpha_{\theta})$$

$$\therefore \Psi = Q = D_\sigma \int_{\phi=0}^{2\pi} \int_{\theta=0}^{\pi} \sigma^2 \sin\theta d\theta d\phi$$

$$Q = D_\sigma 2\pi \int_{\theta=0}^{\pi} \sigma^2 \sin\theta d\theta$$

$$Q = D_\sigma 2\pi \sigma^2 \left[ -\cos\theta \right]_0^{\pi} = D_\sigma 2\pi \sigma^2 [(-1-1)]$$

$$Q = 4\pi \sigma^2 D_\sigma \Rightarrow D_\sigma = \frac{Q}{4\pi \sigma^2}$$

$$\therefore D = D_\sigma \alpha_\sigma \Rightarrow D = \boxed{D = \frac{Q}{4\pi \sigma^2} \alpha_\sigma}$$

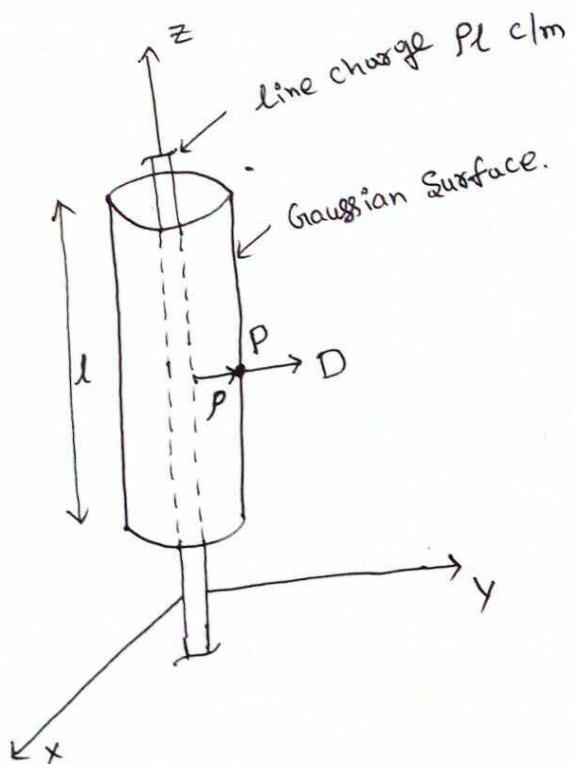
## 2. Infinite Line Charge :

→ Suppose the infinite line of uniform charge  $\rho_L$  c/m lies along the z-axis.

→ To determine  $D$  at a point  $P$ , we choose a cylindrical surface containing  $P$  to satisfy the symmetry condition as shown in figure.

→ The electric flux density  $D$  is constant on and normal to the cylindrical gaussian surface; i.e.  $D = D_p \alpha_p$   
applying Gauss's Law

$$\Psi = Q_{\text{enclosed}} = \oint_S D \cdot dS = D_p \oint_S dS$$



$dS \rightarrow$  Surface area along  $\rho$  direction in cylindrical coordinate system.

$$\therefore \Psi = Q = D_p \int_{z=0}^{+\infty} \int_{\phi=0}^{2\pi} \rho d\phi dz$$

$$Q = D_p \rho l 2\pi = 2\pi \rho l D_p \rightarrow (1)$$

→ Total line charge is  $dQ = \rho_l dl$

$$Q = \int \rho_l dl = \rho_l \cdot l \rightarrow (2)$$

$$\text{Equating eq (1) \& (2)} \quad 2\pi \rho l / D_p = \rho_l \cdot l$$

$$D_p = \frac{\rho_l}{2\pi \rho}$$

$$\therefore D = D_p a_p \Rightarrow D = \frac{\rho_l}{2\pi \rho} a_p$$

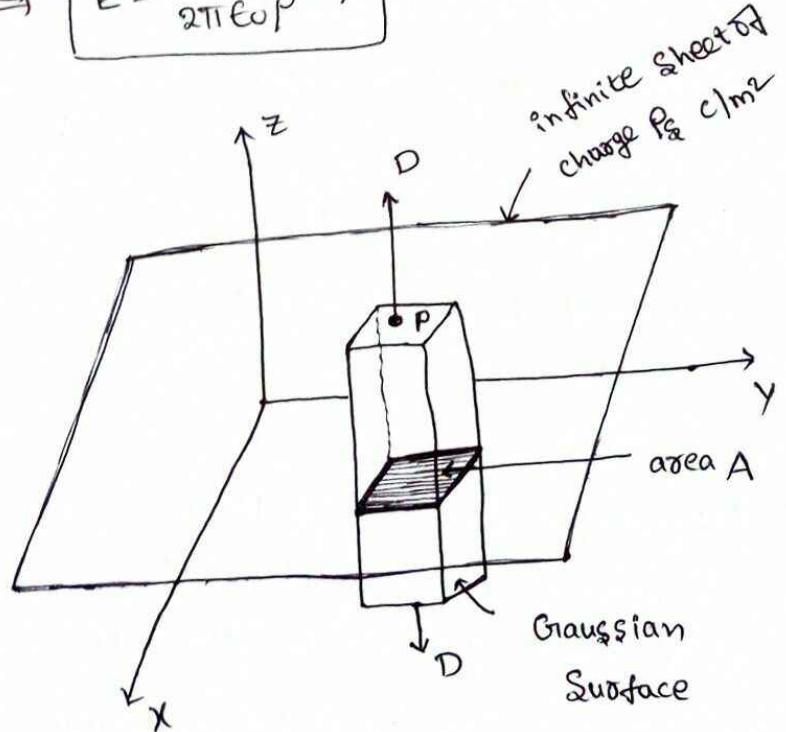
$$\rightarrow \text{Electric field Intensity } E = \frac{D}{\epsilon_0} \Rightarrow E = \frac{\rho_l}{2\pi \epsilon_0 \rho} a_p$$

### 3. Infinite Sheet of charge :

→ Consider the infinite sheet of uniform charge  $\rho_s \text{ C/m}^2$  lying on the  $z=0$  plane.

→ To determine  $D$  at point  $P$ , we choose a rectangular box that is cut symmetrically by the sheet of charge and has two of its faces parallel to the sheet as shown in figure.

→ As  $D$  is normal ( $\perp$  to the sheet i.e.,  $D = D_z a_z$ ) applying Gauss law



$$\Psi = Q_{\text{enclosed}} = \oint_S D_z \cdot dS$$

$$Q = D_z \left[ \int_{\text{top}} dS + \int_{\text{bottom}} dS \right]$$

$$Q = D_z [A + A] = D_z 2A \rightarrow (1)$$

→ Total Surface charge  $dQ = \rho_s dS$

$$Q = \rho_s \int dS = \rho_s A \rightarrow (2)$$

Equating eq (1) & (2)

$$D_z 2A = \rho_s A$$

$$D_z = \frac{\rho_s}{2}$$

$$D = D_z a_z \Rightarrow D = \frac{\rho_s}{2} a_z$$

$$\rightarrow \text{Electric Field Intensity } E = \frac{D}{\epsilon_0} \Rightarrow E = \frac{\rho_s}{2\epsilon_0} a_z$$

$$\int dS = \int dx dy = xy = A$$

$$\begin{aligned} \text{area of rectangle} \\ A = a \times b \end{aligned}$$

#### 4. Uniformly charged Sphere:

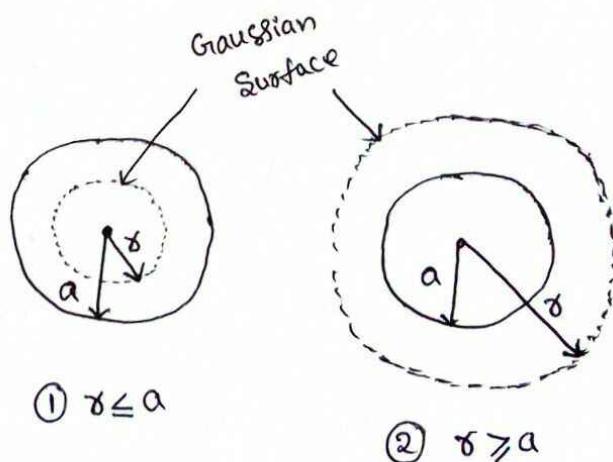
→ Consider a sphere of radius  $a$  with a uniform charge  $\rho_v$   $C/m^3$ .

→ To determine  $D$  everywhere we construct Gaussian Surface for cases  $r \leq a$  &  $r \geq a$  separately.

→ Since the charge has spherical symmetry, it is obvious that a spherical surface is an appropriate Gaussian Surface.

Case ①: for  $r \leq a$ , the total charge enclosed by the spherical surface of radius  $r$ , is shown in figure.

$$Q_{\text{enclosed}} = \int \rho_v dv = \rho_v \int dv$$



$$Q = \rho_V \int_{\phi=0}^{2\pi} \int_{\theta=0}^{\pi} \int_{r=0}^{\infty} r^2 \sin\theta \, dr \, d\theta \, d\phi$$

$$Q = \rho_V 2\pi \left[ -\cos\theta \right]_0^{\pi} \frac{r^3}{3} = \rho_V 2\pi (2) \frac{r^3}{3}$$

$$\therefore Q = \rho_V \frac{4}{3} \pi r^3$$

flux  $\Psi = \oint_D D \cdot dS = D\sigma \oint_S dS$

$$\Psi = D\sigma \int_{\phi=0}^{2\pi} \int_{\theta=0}^{\pi} r^2 \sin\theta \, d\theta \, d\phi$$

$$\Psi = D\sigma 2\pi \left[ -\cos\theta \right]_0^{\pi} r^2 = D\sigma 4\pi r^2 = \Psi$$

$$\therefore Q_{\text{enclosed}} = \Psi \Rightarrow \rho_V \frac{4}{3} \pi r^3 = D\sigma 4\pi r^2 \Rightarrow D\sigma = \frac{\rho_V r}{3}$$

$$\therefore D = D\sigma a\sigma \Rightarrow D = \frac{r}{3} \rho_V a\sigma \quad \rightarrow (1)$$

Case ②: for  $r \geq a$  the ~~total~~ charge enclosed by the surface is entire charge

$$Q = \int \rho_V dv = \rho_V \int dv = \rho_V \int_{\phi=0}^{2\pi} \int_{\theta=0}^{\pi} \int_{r=0}^a r^2 \sin\theta \, dr \, d\theta \, d\phi = \rho_V \frac{4}{3} \pi a^3$$

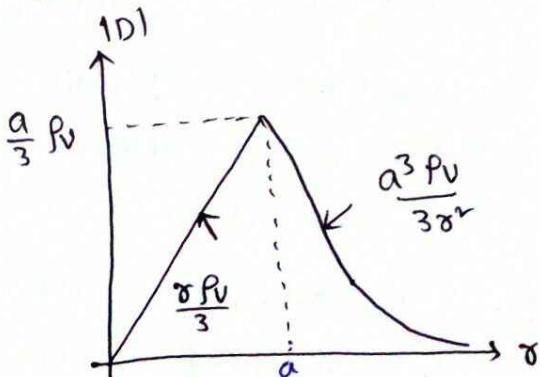
$$\Psi = \oint_D D \cdot dS = D\sigma \int_S dS = D\sigma \int_{\phi=0}^{2\pi} \int_{\theta=0}^{\pi} r^2 \sin\theta \, d\theta \, d\phi = D\sigma 4\pi r^2$$

$$\therefore Q = \Psi \Rightarrow \rho_V \frac{4}{3} \pi a^3 = D\sigma 4\pi r^2 \Rightarrow D\sigma = \frac{a^3}{3r^2} \rho_V$$

$$\therefore D = D\sigma a\sigma \Rightarrow D = \frac{a^3}{3r^2} \rho_V a\sigma \quad \rightarrow (2)$$

→ from eq (1) & (2), D everywhere is given by

$$D = \begin{cases} \frac{r}{3} \rho_V a\sigma & 0 < r \leq a \\ \frac{a^3}{3r^2} \rho_V a\sigma & r \geq a \end{cases}$$



|D| Versus r for a uniform charge sphere.

∴ differential volume  
in spherical coordinate  
system  $\sigma^2 \sin\theta d\theta d\phi$

4.7 Determine  $D$  at  $(4, 0, 3)$  if there is a point charge  $-5\pi \text{ nC}$  at  $(4, 0, 0)$  and a line charge  $3\pi \text{ nC/m}$  along the  $y$ -axis

Sol Electric flux density due to point charge

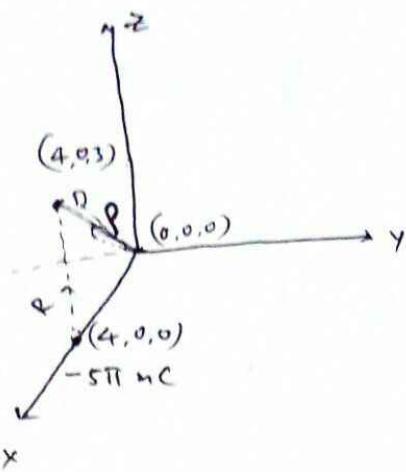
$$D_Q = \epsilon_0 E = \frac{\epsilon_0 Q}{4\pi \epsilon_0 R^2} = \frac{Q}{4\pi R^2} = \frac{Q}{4\pi (R)^3}$$

$$\mathbf{r} = (4ax - 4ax) + (0ay - 0ay) + (3az - 0az)$$

$$R = 0 + 0 + 3az = 3az$$

$$(R) = \sqrt{3^2} = \sqrt{9} = 3$$

$$D = \frac{-5\pi \times 10^{-3}}{4\pi (3)^3} (3az) = \frac{-5 \times 10^{-3} 3az}{4 \times 27} = -0.138 az \text{ nC/m}^2$$



Electric flux density due to line charge

$$D_L = \frac{P_L}{2\pi R^2} az = \frac{P_L}{2\pi (R)^2} P$$

$$P = 4ax + 0 + 3az = 4ax + 3az$$

$$|P| = \sqrt{4^2 + 3^2} = \sqrt{16+9} = \sqrt{25} = 5$$

$$\therefore D_L = \frac{3\pi \times 10^{-3}}{2\pi (5)^2} \times (4ax + 3az) = 0.24ax + 0.18az \text{ nC/m}^2$$

$$\therefore D = D_Q + D_L = 240ax + 42az \text{ nC/m}^2$$

4.8 Given that  $D = z\rho \cos^2 \phi az \text{ C/m}^2$ , calculate the charge density at  $(1, \pi/4, 3)$  and the total charge enclosed by the cylinder of radius 1m ~~with height 2m~~ with  $-2 \leq z \leq 2 \text{ m}$

Sol  $D = z\rho \cos^2 \phi az \text{ C/m}^2$

$$\text{From Gauss law } P_V = \nabla \cdot D = \left( \frac{\partial ax}{\partial x} + \frac{\partial ay}{\partial y} + \frac{\partial az}{\partial z} \right) \cdot (z\rho \cos^2 \phi az) = \frac{\partial z\rho \cos^2 \phi}{\partial z}$$

$$P_V = 1 \rho \cos^2 \phi = \rho \cos^2 \phi ?$$

$$\text{at } (1, \pi/4, 3). \quad P_V = 1 \cdot \cos^2(\pi/4) = 1 \cdot (0.5) = 0.5 \text{ C/m}^2$$

$$\cos(\pi/4) = \frac{1}{\sqrt{2}}$$

$$\begin{aligned} Q &= \int P_V dv = \int \rho \cos^2 \phi \rho d\rho d\phi dz \\ Q &= \int_{z=-2}^2 dz \int_{\phi=0}^{2\pi} \cos^2 \phi d\phi \int_{\rho=0}^1 \rho^2 d\rho = \left[ z \right]_2^2 \left( \frac{\rho^3}{3} \right)_0^{2\pi} \int_{\phi=0}^{2\pi} \left( 1 + \cos^2 \phi \right) d\phi \end{aligned}$$

Value of the cylinder  
 $dv = \rho d\rho d\phi dz$

$$Q = 4 \cdot \frac{1}{3} \cdot \left[ \int_0^{2\pi} \frac{1}{2} d\phi + \int_0^{2\pi} \frac{C \sin \phi}{2} d\phi \right] = \frac{4}{3} \left[ \frac{2\pi}{2} + \frac{(\sin 4\pi - \sin 0)}{2} \right]$$

$$Q = \frac{4\pi}{3} C$$

$\therefore Q = \frac{4\pi}{3} C$

Method 2 from Gauss Law

$$Q = \Psi = \oint D ds = \left[ \int_s + \int_t + \int_b \right] D \cdot ds = \Psi_s + \Psi_t + \Psi_b$$

$$\Psi_s = \int_s D ds = \int_s z p \cos^2 \phi dz \cdot p dp d\phi \Big|_{z=2}$$

Surface of the cylinder  
along z-axis  
 $ds = p dp d\phi$  (top)

$$\Psi_t = 2 \int_{p=0}^1 p^2 dp \int_{\phi=0}^{2\pi} \cos^2 \phi d\phi = 2 \left( \frac{1}{3} \right) \cdot \left( \frac{2\pi}{2} + 0 \right) = \frac{2}{3}\pi$$

$$\therefore \Psi_t = \frac{2\pi}{3}$$

$4 \cdot 12 \times 2 =$   
 $6 \cdot 12$

$$\Psi_b, ds = -p dp d\phi \text{ (bottom)}$$

$$\Psi_b = \int_b D ds = \int_b z p \cos^2 \phi dz \cdot (-p dp d\phi) \Big|_{z=-2}$$

$$\Psi_b = +2 \int_{p=0}^1 p^2 dp \int_{\phi=0}^{2\pi} \cos^2 \phi d\phi = 2 \left( \frac{1}{3} \right) \left( \frac{2\pi}{2} + 0 \right) = \frac{2}{3}\pi$$

$$\therefore \Psi_b = \frac{2\pi}{3}$$

$$\therefore Q = \Psi = \Psi_s + \Psi_t + \Psi_b = 0 + \frac{2\pi}{3} + \frac{2\pi}{3} = \frac{4\pi}{3} C$$

Problem  
(4.5)

The finite sheet  $0 \leq x \leq 1$ ,  $0 \leq y \leq 1$  on the  $z=0$  plane has a charge density  $\rho_s = xy(x^2 + y^2 + 25)^{3/2}$  nC/m<sup>2</sup>. Find

- (a) the total charge on the sheet
- (b) the electric field at  $(0, 0, 5)$ .
- (c) the force experienced by a  $-1$  nC charge located at  $(0, 0, 5)$ .

Sol (a)

$$Q = \int \rho_s dS = \iint_0^1 xy(x^2 + y^2 + 25)^{3/2} dx dy \text{ nC.}$$

$$\text{Let } x^2 = u \Rightarrow \frac{du}{dx} = 2x \Rightarrow x dx = \frac{1}{2} du = \frac{1}{2} d(u^2).$$

$$y^2 = u \Rightarrow \frac{du}{dy} = 2y \Rightarrow y dy = \frac{1}{2} du = \frac{1}{2} d(y^2).$$

Now integrate w.r.t.  $x^2 + y^2$ .

$$\therefore Q = \int_0^1 y \int_0^1 \frac{1}{2} (x^2 + y^2 + 25)^{3/2} d(u^2) dy \text{ nC}$$

$$= \frac{1}{2} \int_0^1 y \left[ \frac{6u^2 + y^2 + 25}{\frac{3}{2} + 1} \right]_0^1 dy$$

$$= \frac{1}{2} \times \frac{2}{5} \int_0^1 \left[ (26 + y^2)^{5/2} - (25 + y^2)^{5/2} \right] \frac{1}{2} d(y^2)$$

$$= \frac{1}{5} \times \frac{1}{2} \left[ \frac{(26 + y^2)^{\frac{5}{2} + 1}}{\frac{5}{2} + 1} - \frac{(25 + y^2)^{\frac{5}{2} + 1}}{\frac{5}{2} + 1} \right]_0^1$$

$$= \frac{1}{10} \times \frac{2}{7} \left[ \left\{ (27)^{7/2} - (26)^{7/2} \right\} - \left\{ (25)^{7/2} - (25)^{7/2} \right\} \right]$$

$$= \frac{1}{35} \left[ (27)^{7/2} + (25)^{7/2} - 2(26)^{7/2} \right]$$

$$Q = 33.15 \text{ nC.}$$

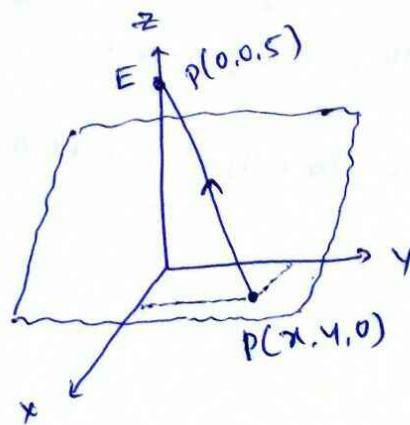
(b)

$$E = \int \frac{\rho_s dS}{4\pi\epsilon_0 R^2} dR$$

$$R = (0-x)\hat{x} + (0-y)\hat{y} + (5-0)\hat{z}$$

$$R = -x\hat{x} - y\hat{y} + 5\hat{z}$$

$$|R| = \sqrt{x^2 + y^2 + 25}$$



$$E = \int \frac{Ps \, dS}{4\pi \epsilon_0 R^3} = \int_0^1 \int_0^1 \frac{\pi y(x^2 + y^2 + z^2)^{-1/2} \times 10^{-9} (-xax - yay + za^z)}{4\pi \times \frac{1}{36\pi \times 10^9} (x^2 + y^2 + z^2)^{3/2}} \, dx \, dy$$

$$E = 9 \times 10^{-9} \left[ - \int x^2 \, dx \int y \, dy \, ax - \int x \, dx \int y^2 \, dy \, ay + \int z \, dz \int y \, dy \, az \right]$$

$$E = 9 \left[ - \left[ \frac{x^3}{3} \right]_0^1 \left[ \frac{y^2}{2} \right]_0^1 ax - \left[ \frac{x^2}{2} \right]_0^1 \left[ \frac{y^3}{3} \right]_0^1 ay + 5 \left[ \frac{z^2}{2} \right]_0^1 \left[ \frac{y^2}{2} \right]_0^1 az \right]$$

$$E = 9 \left[ - \frac{1}{3} \times \frac{1}{2} ax - \frac{1}{2} \times \frac{1}{3} ay + 5 \times \frac{1}{2} \times \frac{1}{2} az \right] = 9 \left[ -\frac{1}{6} ax - \frac{1}{6} ay + \frac{5}{4} az \right]$$

$$\boxed{E = -1.5ax - 1.5ay + 11.25az \text{ V/m (in N/C)}}$$

(e)  $F = EQ = (-1 \times 10^{-3}) (-1.5ax - 1.5ay + 11.25az)$

$$\boxed{F = 1.5ax + 1.5ay - 11.25az \text{ mN.}}$$

Problem  
PE 4.7

A point charge  $Q = 30\text{nC}$  is located at the origin while plane  $y=3$  carries charge  $10\text{nC/m}^2$ . Find  $D$  at  $(0, 4, 3)$

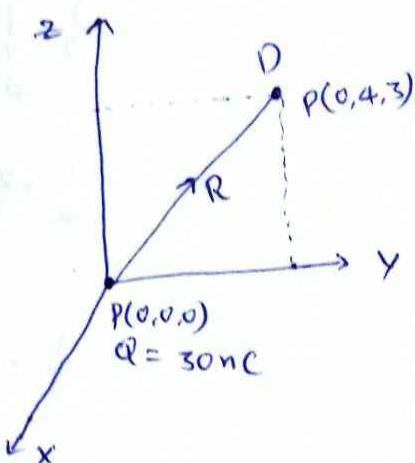
Sol

$$R = (0-0)ax + (4-0)ay + (3-0)az = 4ay + 3az$$

$$|R| = \sqrt{16+9} = 5$$

$$D_a = \frac{Q}{4\pi |R|^3} R = \frac{30 \times 10^{-9} (4ay + 3az)}{4\pi \times 125}$$

$$D_a = 0.0763ay + 0.0573az \text{ nC/m}^2$$



→ plane  $y=3$  means sheet is in the  $xz$  plane. Hence  $P_s$  is along  $ay$  direction

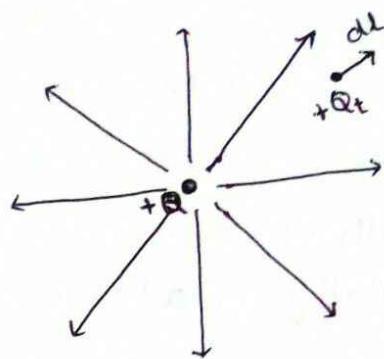
$$\therefore D_s = \frac{P_s}{2} ay = \frac{10 \times 10^{-9}}{2} ay = 5 \times 10^{-9} ay$$

$$D_s = 5ay \text{ nC/m}^2$$

$$\boxed{\therefore D = D_a + D_s = 5.076ay + 0.0573az \text{ nC/m}^2}$$

Work Done :

Consider a positive charge  $Q$  and its electric field  $E$ . If a positive test charge  $q_t$  is placed in this field, it will move due to force of repulsion. Let the movement of the charge  $q_t$  is  $dl$ .



- To keep the charge in equilibrium, it is necessary to apply the force which is equal and opposite to the force exerted by the field in the direction  $dl$ .
- Thus keeping the charge in equilibrium means we are moving a charge  $q_t$  through the distance  $dl$  in opposite direction to that of field  $E$ . Hence the work is done.

From coulomb's law, the force on  $q_t$  is  $F = E q_t$ . So that the work done in displacing the charge by  $dl$  is

$$dw = -F \cdot dl = -E q_t \cdot dl$$

The negative sign indicates that the work is being done by an external agent. Thus the total work done is

$$W = -q_t \int E \cdot dl$$

Electric Potential :

The work done in moving a unit charge from one point to another point in an electric field  $E$  is called electric potential (or) Potential difference (or) Electric scalar Potential (or) Potential.

- It is denoted by  $V$  (or)  $V_{AB}$  and measured in volts. It is a scalar quantity.

$$\therefore V = \frac{W}{Q} = - \int E \cdot dl$$

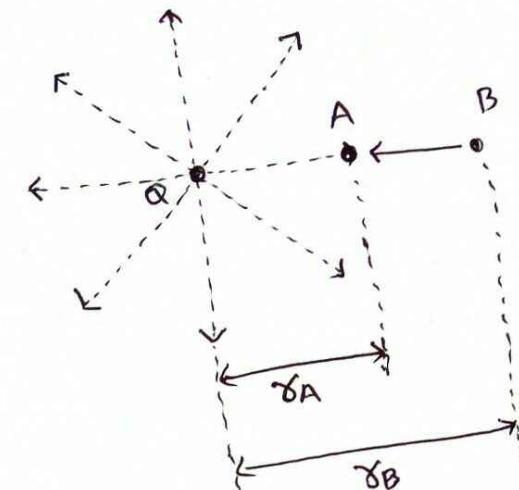
$$\therefore V = - \int E \cdot dl$$

Here  $Q$  is selected as unit test charge  $q_t$  from work done.

## Electric Potential Due to Point charge :

Consider a point charge  $Q$ , located at the origin of a spherical coordinate system, producing  $E$  radially in all the directions as shown in the figure.

Assuming free space, the field  $E$  due to a point charge  $Q$  at a point having radial distance  $\sigma$  from the origin is given by  $E = \frac{Q}{4\pi\epsilon_0\sigma^2} a_\sigma$



$$V = - \int E \cdot d\ell = - \int_{r_B}^{r_A} \frac{Q}{4\pi\epsilon_0\sigma^2} a_\sigma \cdot d\sigma$$

$$V = \frac{-Q}{4\pi\epsilon_0} \int \frac{1}{\sigma^2} d\sigma = \frac{-Q}{4\pi\epsilon_0} \left(-\frac{1}{\sigma}\right)$$

$$V = \frac{Q}{4\pi\epsilon_0 r}$$

- Consider a unit charge which is placed at Point B which is at a radial distance  $r_B$  from the origin. It is moved against the direction of  $E$  from Point B to Point A. Point A is at a radial distance  $r_A$  from the origin.
- The differential length in spherical system is

$$d\ell = dr a_\sigma + \sigma d\theta a_\theta + \sigma \sin\theta d\phi a_\phi$$

Hence the potential difference  $V_{AB}$  b/w Points A & B is given by

$$V_{AB} = - \int_B^A E \cdot d\ell \quad \text{But } B \Rightarrow r_B \quad \text{and} \quad A \Rightarrow r_A$$

$$\therefore V_{AB} = - \int_{r_B}^{r_A} \left( \frac{Q}{4\pi\epsilon_0\sigma^2} a_\sigma \right) \cdot (dr a_\sigma) = - \int_{r_B}^{r_A} \frac{Q}{4\pi\epsilon_0\sigma^2} d\sigma$$

$$\therefore V_{AB} = \frac{-Q}{4\pi\epsilon_0} \int_{r_B}^{r_A} \sigma^{-2} d\sigma = \frac{-Q}{4\pi\epsilon_0} \left[ -\frac{1}{\sigma} \right]_{r_B}^{r_A} = \frac{-Q}{4\pi\epsilon_0} \left[ -\frac{1}{r_A} + \frac{1}{r_B} \right]$$

$$\therefore V_{AB} = \frac{Q}{4\pi\epsilon_0} \left[ \frac{1}{r_A} - \frac{1}{r_B} \right] V$$

When  $r_B > r_A$ ,  $\frac{1}{r_B} < \frac{1}{r_A}$  and  $V_{AB}$  is positive. This indicates that the work is done by external source in moving unit charge from B to A.

Absolute Potential :

The absolute Potential at any Point in an electric field is defined as the work done in moving a unit test charge from infinity to the point, against the direction of the field.

- Consider potential difference  $V_{AB}$  due to movement of unit charge from B to A in a field of Point charge  $Q$  is

$$V_{AB} = \frac{Q}{4\pi\epsilon_0} \left[ \frac{1}{r_A} - \frac{1}{r_B} \right]$$

Now charge is moved from infinity to Point A i.e  $r_B = \infty \Rightarrow \frac{1}{r_B} = \frac{1}{\infty} = 0$

$$\therefore V_{AB} = \frac{Q}{4\pi\epsilon_0} \left[ \frac{1}{r_A} - 0 \right] = \frac{Q}{4\pi\epsilon_0 r_A} \text{ V.}$$

It is also called potential of Point A denoted by  $V_A$

$$\boxed{\therefore V_A = \frac{Q}{4\pi\epsilon_0 r_A} \text{ V.}}$$

This is also called absolute Potential of Point A.

Similarly absolute Potential of Point B can be defined as

$$\boxed{\therefore V_B = \frac{Q}{4\pi\epsilon_0 r_B} \text{ V}}$$

This is the work done in moving unit charge from infinity at Point B.

- Hence the Potential difference can be expressed as the difference between the absolute Potentials of the two points.

$$\therefore V_{AB} = V_A - V_B$$

- Therefore the potential (or) absolute Potential at any Point which is at a distance  $r$  from the origin of a spherical system, where Point charge  $Q$  is located, is given by

$$\boxed{V = \frac{Q}{4\pi\epsilon_0 r}} \longrightarrow (1)$$

→ If the Point charge  $Q$  in eq(1) is not located at the origin but at Point whose Position vector  $\vec{r}'$ , the Potential  $V(x,y,z)$  at  $\vec{r}$  becomes

$$V(x,y,z) = V(\vec{r}) = \frac{Q}{4\pi\epsilon_0 |\vec{r} - \vec{r}'|}$$

→ For  $n$  point charges  $Q_1, Q_2, \dots, Q_n$  located at Points with position vectors  $\vec{r}_1, \vec{r}_2, \dots, \vec{r}_n$ , the Potential at  $\vec{r}$  is

$$V(\vec{r}) = \frac{Q_1}{4\pi\epsilon_0 |\vec{r} - \vec{r}_1|} + \frac{Q_2}{4\pi\epsilon_0 |\vec{r} - \vec{r}_2|} + \dots + \frac{Q_n}{4\pi\epsilon_0 |\vec{r} - \vec{r}_n|}$$

$$V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \sum_{k=1}^n \frac{Q_k}{|\vec{r} - \vec{r}_k|} \rightarrow (2) \text{ (Point charge)}$$

→ For continuous charge distributions, we replace  $Q_k$  in eq(2) with charge element  $P_l dl$ ,  $P_s ds$ ,  $P_v dv$  and the summation becomes an integration, so the Potential at  $\vec{r}$  becomes

$$V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \int_L \frac{P_l(\vec{r}') dl}{|\vec{r} - \vec{r}'|} \quad (\text{line charge})$$

$$V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \int_S \frac{P_s(\vec{r}') dS}{|\vec{r} - \vec{r}'|} \quad (\text{surface charge})$$

$$V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \int_V \frac{P_v(\vec{r}') dv}{|\vec{r} - \vec{r}'|} \quad (\text{volume charge})$$

### Problem

TWO Point charges  $-4 \mu C$  and  $5 \mu C$  are located at  $(2, -1, 3)$  and  $(0, 4, -2)$ , respectively. Find the Potential at  $(1, 0, 1)$ , assuming zero potential at infinity.

Sol

Given  $Q_1 = -4 \mu C$ ,  $Q_2 = 5 \mu C$

$$\text{Potential } V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \left[ \frac{Q_1}{|\vec{r} - \vec{r}_1|} + \frac{Q_2}{|\vec{r} - \vec{r}_2|} \right]$$

$$|\vec{r} - \vec{r}_1| = |(1, 0, 1) - (2, -1, 3)| = |-1, 1, -2| = \sqrt{1+1+4} = \sqrt{6}$$

$$|\vec{r} - \vec{r}_2| = |(1, 0, 1) - (0, 4, -2)| = |1, -4, 3| = \sqrt{1+16+9} = \sqrt{26}$$

$$\therefore V(1, 0, 1) = \frac{10^{-6}}{4\pi \times \frac{10^{-9}}{36\pi}} \left[ \frac{-4}{\sqrt{6}} + \frac{5}{\sqrt{26}} \right]$$

$$= 9 \times 10^3 (-1.633 + 0.9806)$$

$$= -5.872 \text{ kV}$$

W.Hart Given the electric flux density  $D = 0.3 \text{ a} \text{ s nC/m}^2$  in free space.

(a) Find electric field  $E$  at a point  $P(\gamma=2, \theta=25^\circ, \phi=90^\circ)$

(b) Find the total charge within radius of sphere  $\gamma_1=3$

(c) Find the total electric flux leaving the surface of sphere  $\gamma=4$ .

Sol (a) Given  $D = 0.3 \text{ a} \text{ s nC/m}^2$

$$\eta = \epsilon_0 E \Rightarrow E = \frac{D}{\epsilon_0} = 36\pi \times 10^9 \times 0.3 \times 10^{-9} \text{ a} \text{ s}$$

$$E = 67.8 \text{ a} \text{ s V/m}$$

If  $D = 0.3 \text{ a} \text{ s nC/m}^2$

then (a) 135.5 a s V/m

(b) 305 nC

(c) 965 nC

(b)  $Q(\gamma=3)$

$$Q = \int_S D \cdot dS = \int (0.3 \text{ a} \text{ s}) a \gamma \cdot (\gamma^2 \sin \theta d\theta d\phi) a \gamma$$

$$Q = \int_{\phi=0}^{2\pi} \int_{\theta=0}^{\pi} (0.3 \gamma^3) \sin \theta d\theta d\phi \Big|_{\gamma=3}$$

$$Q = (0.3) (3)^3 2\pi (-\cos \theta)_0^{\pi} = (0.3) (2\pi) (2\pi) (2)$$

$$Q = 101.78 \text{ nC}$$

$$\text{In. } kT = D = \frac{Q}{4\pi r^2} \Rightarrow Q = D 4\pi r^2$$

(c)  $\Psi = Q (\gamma=4)$

$$\Psi = D (4\pi \gamma^2) = 0.3 \times (4\pi \times 4^2) = (0.3)(4\pi)(4^3)$$

$$\Psi = (0.3)(4\pi)(4)^3 = 241.27 \text{ nC}$$

Given the Potential  $V = \frac{10}{r^2} \sin\theta \cos\phi$ ,

(a) find the electric flux density  $D$  at  $(2, \pi/2, 0)$ .

(b) Calculate the work done in moving a  $10 \mu C$  charge from point A( $1, 30^\circ, 120^\circ$ ) to B( $4, 90^\circ, 60^\circ$ ).

Sol

$$(a) D = \epsilon_0 E$$

$$E = -\nabla V = - \left[ \frac{\partial V}{\partial r} \hat{r} + \frac{1}{r} \frac{\partial V}{\partial \theta} \hat{\theta} + \frac{1}{r \sin\theta} \frac{\partial V}{\partial \phi} \hat{\phi} \right]$$

$$E = - \left[ 10 \sin\theta \cos\phi \frac{d}{dr} \hat{r} + \frac{1}{r} \frac{10}{r^2} \cos\phi \frac{d}{d\theta} \hat{\theta} + \frac{1}{r \sin\theta} \frac{10}{r^2} \frac{\sin\theta}{\cos\phi} \frac{d}{d\phi} \hat{\phi} \right]$$

$$E = - \left[ 10 \sin\theta \cos\phi \frac{2}{r^3} \hat{r} + \frac{10}{r^3} \cos\phi \cos\theta \hat{\theta} + \frac{10}{r^3} (-\sin\phi) \hat{\phi} \right]$$

$$E = \frac{20}{r^3} \sin\theta \cos\phi \hat{r} - \frac{10}{r^3} \cos\theta \cos\phi \hat{\theta} + \frac{10}{r^3} \sin\phi \hat{\phi}$$

$$\text{At } (2, \pi/2, 0) \quad D = \epsilon_0 E \quad (r=2, \theta=\pi/2, \phi=0)$$

$$D = \epsilon_0 \left[ \frac{20}{8} (1) \hat{r} - 0 \hat{\theta} + 0 \hat{\phi} \right]$$

$$D = \epsilon_0 (2.5) \hat{r} \text{ C/m}^2$$

$$D = 22.1 \hat{r} \text{ pC/m}^2$$

$$(b) W = -Q \int_A^B \mathbf{E} \cdot d\mathbf{l} = Q V_{AB} = Q (V_B - V_A)$$

$$W = 10 \left[ \frac{10}{16} \sin 90^\circ \cos 60^\circ - \frac{10}{1} \sin 30^\circ \cos 120^\circ \right] \cdot 10^6$$

$$W = 10 \left[ \frac{10}{32} - \frac{-5}{2} \right] \cdot 10^6 = 28.125 \text{ uJ}$$

$$\therefore W = 28.125 \text{ MJ}$$

## Relation between E and V (Potential Gradient) :

Electric field intensity (E) is equal to the negative gradient of electric potential (V).

$$\text{i.e. } \boxed{E = -\nabla V}$$

Proof:

The electric Potential is given by  $V = - \int E \cdot d\ell$

Apply vector differential operators on both sides

$$\nabla V = -\nabla \int E \cdot d\ell$$

$$-\nabla V = \left( \frac{\partial}{\partial x} ax + \frac{\partial}{\partial y} ay + \frac{\partial}{\partial z} az \right) \int (Exax + Eyah + Ezaz) \cdot (dxax + dyay + dzaz)$$

$$-\nabla V = \left( \frac{\partial}{\partial x} ax + \frac{\partial}{\partial y} ay + \frac{\partial}{\partial z} az \right) \int (Exdx + Eyd\gamma + Ezdz)$$

$$-\nabla V = \frac{\partial}{\partial x} \int Ex dx ax + \frac{\partial}{\partial y} \int Ey dy ay + \frac{\partial}{\partial z} \int Ez dz az$$

$$-\nabla V = Exax + Eyah + Ezaz$$

[: Integration & Differentiation  
are cancelled ]

$$-\nabla V = E$$

$$\therefore \boxed{E = -\nabla V}$$

(Q8)

$$\text{The Electric field intensity } E = \frac{\alpha}{4\pi\epsilon_0 r^2} ax \quad \rightarrow (1)$$

$$\text{we know that } \nabla \left( \frac{1}{r} \right) = -\frac{1}{r^2} ar \Rightarrow \frac{1}{r^2} ar = -\nabla \left( \frac{1}{r} \right) \quad \rightarrow (2)$$

$$\text{Substitute (2) in (1)} \quad E = \frac{\alpha}{4\pi\epsilon_0} \left[ -\nabla \left( \frac{1}{r} \right) \right] = -\nabla \frac{\alpha}{4\pi\epsilon_0 r} = -\nabla V \quad | : V = \frac{\alpha}{4\pi\epsilon_0 r}$$

$$\therefore \boxed{E = -\nabla V} \quad \text{Electric field is negative potential gradient.}$$

\* Potential is scalar. Gradient of scalar function is a vector function.

## Maxwell's Two Equations for Electro Static field :

→ Gauss's law in Point form is called Maxwell's first equation.

Equation 1 : The divergence of the electric flux density ( $D$ ) is equal to the volume charge density ( $P_V$ ).

i.e., 
$$\nabla \cdot D = P_V$$

Proof : from the definition of Gauss's law

$$Q = \Psi = \oint_S D \cdot dS$$

Apply divergence theorem to the above equation.

$$Q = \oint_S D \cdot dS = \int_V \nabla \cdot D \, dv \quad \rightarrow (1)$$

→ Total charge enclosed by the volume is

$$Q = \int_V P_V \, dv \quad \rightarrow (2)$$

equating eq(1) & eq(2)

$$\int_V \nabla \cdot D \, dv = \int_V P_V \, dv$$

$$\boxed{\nabla \cdot D = P_V}$$

⇒ Conservative field (or) irrotational field is called Maxwell's second equation.

Equation 2 :

In conservative field the curl of electric field intensity is equal to zero

i.e., 
$$\nabla \times E = 0$$

Proof :

- The line integral of a vector field  $\mathbf{E}$  around a closed path is zero, such field is called conservative field.
- Physically no work is done in moving a charge along a closed path in an electrostatic field.
- therefore within the closed loop, potential is zero. i.e.,  $V_{AB} + V_{BA} = 0$   
 $\Rightarrow V = 0$

$$\therefore - \oint \mathbf{E} \cdot d\mathbf{l} = 0 \Rightarrow \oint \mathbf{E} \cdot d\mathbf{l} = 0$$

Apply Stokes theorem to the above equation.

$$\oint \mathbf{E} \cdot d\mathbf{l} = \int_S (\nabla \times \mathbf{E}) \cdot d\mathbf{s} = 0$$

$$\therefore \nabla \times \mathbf{E} = 0$$

### Poisson's and Laplace's Equations :

Poisson's Equation :

$$\nabla \cdot \mathbf{D} = \nabla^2 V = -\frac{\rho}{\epsilon_0}$$

Proof :

from Maxwell's first equation

$$\nabla \cdot \mathbf{D} = \rho_V$$

$$\nabla \cdot \epsilon_0 \mathbf{E} = \rho_V$$

$$\epsilon_0 \nabla \cdot \mathbf{E} = \rho_V$$

$$\epsilon_0 \nabla \cdot (-\nabla V) = \rho_V$$

$$-\nabla^2 V = \frac{\rho_V}{\epsilon_0}$$

$$\therefore \mathbf{D} = \epsilon_0 \mathbf{E}$$

$$\therefore \mathbf{E} = -\nabla V$$

$$\therefore \nabla \cdot \mathbf{D} = \nabla^2 V = -\frac{\rho_V}{\epsilon_0}$$

This equation is called Poisson's equation

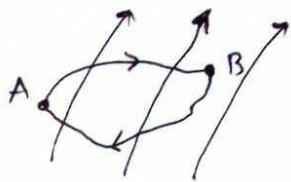


fig: Conservative nature  
of electrostatic field.

## Laplace's Equation :

$$\boxed{\nabla^2 V = 0}$$

Proof :

Consider a charge free region (insulator). The value of  $\rho$  is zero since there are no free charges in dielectrics.  
Substitute  $\rho=0$  in Poisson's equation.

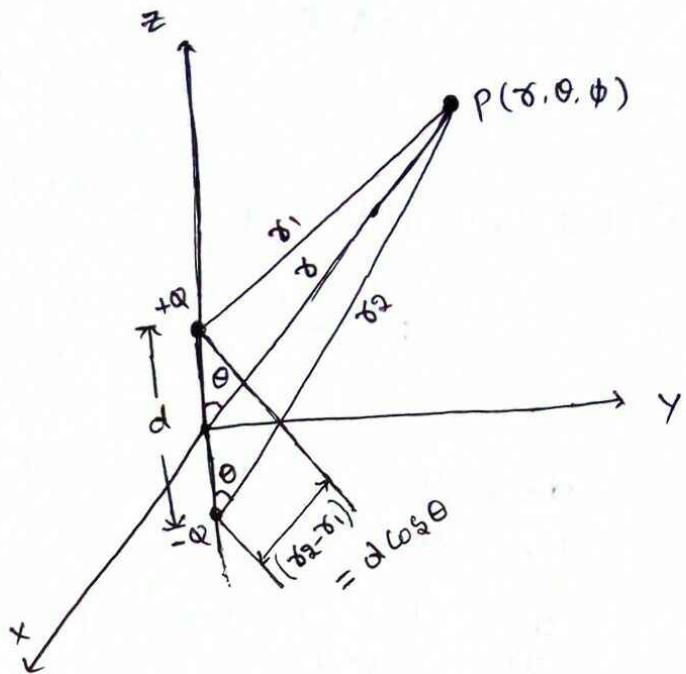
$$\nabla^2 V = \frac{-\rho}{\epsilon_0} = 0$$

$\therefore \nabla^2 V = 0$  This equation is called Laplace's equation.

## Dipole and Dipole Moment :

### Electric Dipole :

→ An electric dipole is formed when two point charges of equal magnitude but opposite sign are separated by a small distance.



- Consider a point  $P(r, \theta, \phi)$  in a spherical coordinate system. It is required to find  $E$  due to an electric dipole at point P.
- The distance of separation of charges i.e.,  $d$  is very small compared to the distances  $r_1$ ,  $r_2$  and  $r$ .

\* In spherical coordinates, the potential at point P due to the charge  $+Q$  is given by

$$V_1 = \frac{+Q}{4\pi\epsilon_0 r_1}$$

\* The potential at Point P due to the charge  $-Q$  is given by

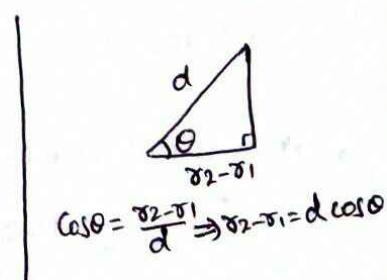
$$V_2 = \frac{-Q}{4\pi\epsilon_0 r_2}$$

\* The total Potential at Point P is the algebraic sum of  $V_1$  and  $V_2$

$$\therefore V_1 + V_2 = V = \frac{Q}{4\pi\epsilon_0 r_1} - \frac{Q}{4\pi\epsilon_0 r_2}$$

$$\therefore V = \frac{Q}{4\pi\epsilon_0} \left[ \frac{1}{r_1} - \frac{1}{r_2} \right]$$

$$V = \frac{Q}{4\pi\epsilon_0} \left[ \frac{r_2 - r_1}{r_1 r_2} \right]$$



$$\cos\theta = \frac{r_2 - r_1}{d} \Rightarrow r_2 - r_1 = d \cos\theta$$

→ where  $r_1$  &  $r_2$  are the distances between P &  $+Q$  and P &  $-Q$  respectively.

If  $r \gg d$ ,  $r_2 - r_1 \approx d \cos\theta$ ,  $r_1 r_2 \approx r^2$ . Then above equation becomes

$$V = \frac{\alpha}{4\pi\epsilon_0} \left[ \frac{d \cos\theta}{r^2} \right]$$

This is the potential

∴ Electric field  $E = -\nabla V$

$$E = - \left[ \frac{\partial V}{\partial r} \hat{a}_r + \frac{1}{r} \frac{\partial V}{\partial \theta} \hat{a}_\theta + \frac{1}{r \sin\theta} \frac{\partial V}{\partial \phi} \hat{a}_\phi \right] \rightarrow (1)$$

$$\frac{\partial V}{\partial r} = \frac{\partial}{\partial r} \left[ \frac{\alpha}{4\pi\epsilon_0} \frac{d \cos\theta}{r^2} \right] = \frac{\alpha}{4\pi\epsilon_0} \frac{d \cos\theta}{r^2} \frac{d}{dr} (r^{-2})$$

$$\frac{\partial V}{\partial r} = \frac{\alpha d \cos\theta}{4\pi\epsilon_0} (-2 \cdot r^{-2-1}) = -\frac{\alpha d \cos\theta}{4\pi\epsilon_0} \left( \frac{2}{r^3} \right) = -\frac{2\alpha d \cos\theta}{4\pi\epsilon_0 r^3}$$

$$\frac{\partial V}{\partial \theta} = \frac{\partial}{\partial \theta} \left[ \frac{\alpha}{4\pi\epsilon_0} \frac{d \cos\theta}{r^2} \right] = \frac{\alpha d}{4\pi\epsilon_0 r^2} (-\sin\theta) = -\frac{\alpha d \sin\theta}{4\pi\epsilon_0 r^2}$$

$$\frac{\partial V}{\partial \phi} = 0$$

Substitute (2) in (1)

$$E = - \left[ -\frac{2\alpha d \cos\theta}{24\pi\epsilon_0 r^3} \hat{a}_r + \frac{1}{r} \left( -\frac{\alpha d \sin\theta}{4\pi\epsilon_0 r^2} \right) \hat{a}_\theta + 0 \right]$$

$$E = \frac{\alpha d}{4\pi\epsilon_0 r^3} [2 \cos\theta \hat{a}_r + \sin\theta \hat{a}_\theta]$$

Dipole moment: It is the product of charge and the separation between charges. It is a vector quantity & has the direction along distance vector  $\vec{d}$  i.e.  $\vec{P} = \alpha \vec{d}$

$$\text{since } d \cos\theta = d \cos\theta = d \cdot \hat{a}_r$$

\* Then dipole moment can be written as

$$V = \frac{\alpha}{4\pi\epsilon_0} \left( \frac{d \cdot \hat{a}_r}{r^2} \right) = \frac{\alpha d \cdot \hat{a}_r}{4\pi\epsilon_0 r^2}$$

$$\therefore V = \frac{P \cdot \hat{a}_r}{4\pi\epsilon_0 r^3} \quad \therefore \hat{a}_r = \frac{\vec{r}}{|\vec{r}|}$$

\* If the dipole centre is not at the origin but at  $r'$ , then dipole moment

can be written as

$$V(r) = \frac{P \cdot (\vec{r} - \vec{r}')}{4\pi\epsilon_0 |\vec{r} - \vec{r}'|^3}$$

## Energy Density in Electrostatic Fields :

- To determine the energy present in an assembly of charges, first determine the amount of work necessary to assemble them.
- Suppose we wish to position consider three point charges  $Q_1, Q_2$  and  $Q_3$  initially in an empty space shown in fig.

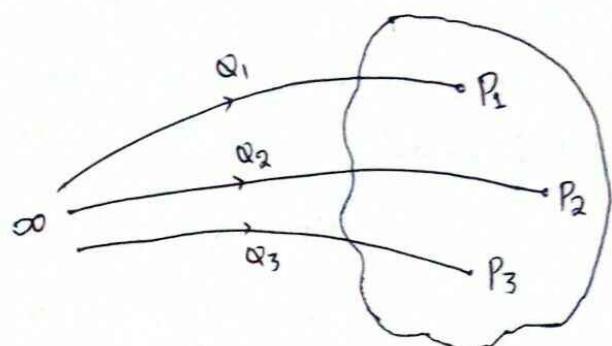


fig: Assembling of charges

- No work is required to transfer  $Q_1$  from infinity to  $P_1$  because the space is initially charge free and there is no electric field [ $w=0$ ].
- The work done in transferring  $Q_2$  from infinity to  $P_2$  is equal to the product of  $Q_2$  and the potential  $V_{21}$  at  $P_2$  due to  $Q_1$ . [i.e.,  $Q_2 V_{21}$ ].
- Similarly, the work done in transferring  $Q_3$  from infinity to  $P_3$  is equal to the  $Q_3 (V_{32} + V_{31})$ , where  $V_{32}$  &  $V_{31}$  are the potentials at  $P_3$  due to  $Q_2$  &  $Q_1$  respectively.

\* Hence the total work done in positioning the three charges is

$$\begin{aligned} WE &= W_1 + W_2 + W_3 \\ &= 0 + Q_2 V_{21} + Q_3 (V_{31} + V_{32}) \end{aligned} \quad \longrightarrow (1)$$

\* If the charges were positioned in reverse order,

$$\begin{aligned} WE &= W_3 + W_2 + W_1 \\ &= 0 + Q_2 V_{23} + Q_1 (V_{12} + V_{13}) \end{aligned} \quad \longrightarrow (2)$$

- where  $V_{23}$  is the potential at  $P_2$  due to  $Q_3$ ,  $V_{12}$  &  $V_{13}$  are the potential at  $P_1$  due to  $Q_2$  &  $Q_3$ . Adding eq (1) & (2)

$$2WE = Q_1(V_{12} + V_{13}) + Q_2(V_{21} + V_{23}) + Q_3(V_{31} + V_{32})$$

$$2WE = Q_1 V_1 + Q_2 V_2 + Q_3 V_3$$

$$WE = \frac{1}{2} (Q_1 V_1 + Q_2 V_2 + Q_3 V_3)$$

where  $V_1, V_2$  &  $V_3$  are total potentials at  $P_1, P_2$  &  $P_3$  respectively.

→ If there are n-point charges, the above equation becomes

$$W_E = \frac{1}{2} \sum_{k=1}^n Q_k V_k$$

Joules.

→ If instead of point charges, the region has a continuous charge distribution, the summation becomes integration i.e.,

$$W_E = \frac{1}{2} \int \rho_l V dl \rightarrow \text{line charge}$$

$$W_E = \frac{1}{2} \int \rho_s V dS \rightarrow \text{Surface charge}$$

$$W_E = \frac{1}{2} \int \rho_v V dv \rightarrow \text{Volume charge}$$

→ Since  $\rho_v = \nabla \cdot D$  then volume charge equation becomes

$$W_E = \frac{1}{2} \int (\nabla \cdot D) V dv$$

→ But for any vector A & scalar V, the identity  $(\nabla \cdot A)V = \nabla \cdot VA - A \cdot \nabla V$

Apply this identity to the above equation.

$$W_E = \frac{1}{2} \int (\nabla \cdot VD) dv - \frac{1}{2} \int (D \cdot \nabla V) dv$$

Apply divergence theorem to the 1<sup>st</sup> term on the right-hand side of eq.

$$W_E = \frac{1}{2} \oint_S (V D) \cdot dS - \frac{1}{2} \int_V (D \cdot \nabla V) dv$$

The work done for a closed loop is zero.

$$W_E = 0 - \frac{1}{2} \int_V (D \cdot \nabla V) dv$$

$$W_E = \frac{1}{2} \int_V (D \cdot E) dv = \frac{1}{2} \int_V (\epsilon_0 E \cdot E) dv$$

$$\begin{aligned} \therefore E &= -\nabla V \\ D &= \epsilon_0 E \end{aligned}$$

$$\therefore W_E = \frac{1}{2} \int_V (D \cdot E) dv = \frac{1}{2} \int_V \epsilon_0 E^2 dv = \frac{1}{2} \int_V \frac{D^2}{\epsilon_0} dv$$

→ from this, we can define electrostatic energy density  $w_E$  (in J/m<sup>3</sup>) as

$$w_E = \frac{dW_E}{dv} = \frac{1}{2} D \cdot E = \frac{1}{2} \epsilon_0 E^2 = \frac{1}{2} \frac{D^2}{\epsilon_0}$$

$$\Rightarrow W_E = \int w_E dv$$

Problem  
4.13)

TWO dipoles with dipole moments  $-5az \text{ nC/m}$  and  $9az \text{ nC/m}$  are located at Points  $(0, 0, -2)$  &  $(0, 0, 3)$  respectively. Find Potential at origin.

S<sub>c</sub>

$$\text{Dipole moment for two point } V = \frac{1}{4\pi\epsilon_0} \left[ \frac{p_1 \cdot r_1}{|r_1|^3} + \frac{p_2 \cdot r_2}{|r_2|^3} \right]$$

$$\therefore p_1 = -5az \text{ & } r_1 = (0, 0, 0) - (0, 0, -2) = 2az, |r_1| = \sqrt{2^2} = 2.$$

$$p_2 = 9az \text{ & } r_2 = (0, 0, 0) - (0, 0, 3) = -3az, |r_2| = \sqrt{3^2} = 3.$$

$$V = \frac{1}{4\pi\epsilon_0} \left[ \frac{-5 \cdot (2)}{2^3} + \frac{9 \cdot (-3)}{3^3} \right] \times 10^{-9}$$

$$V = 9 \times 10^9 \left[ \frac{-10}{8} - \frac{27}{27} \right] \times 10^{-9}$$

$$V = 9 \left( -\frac{10}{8} - 1 \right) = 9 \left( -\frac{10+8}{8} \right) = 9 \left( -\frac{18}{8} \right) = -20.25 \text{ V}$$

$$\therefore V = -20.25$$

Problem  
3(a)

The point charges  $-1 \text{ nC}$ ,  $4 \text{ nC}$  and  $3 \text{ nC}$  are located at  $(0, 0, 0)$ ,  $(0, 0, 1)$  and  $(1, 0, 0)$  respectively. Find the energy in the system

S<sub>c</sub>

$$Q_1 = -1 \text{ nC at } (0, 0, 0)$$

$$Q_2 = 4 \text{ nC at } (0, 0, 1)$$

$$Q_3 = 3 \text{ nC at } (1, 0, 0)$$

$$W = W_1 + W_2 + W_3 = 0 + Q_2 V_{21} + Q_3 (V_{31} + V_{32})$$

$$W = Q_2 \left[ \frac{Q_1}{4\pi\epsilon_0 |(0, 0, 1) - (0, 0, 0)|} \right] + \frac{Q_3}{4\pi\epsilon_0} \left[ \frac{Q_1}{|(1, 0, 0) - (0, 0, 0)|} + \frac{Q_2}{|(1, 0, 0) - (0, 0, 1)|} \right]$$

$$W = Q_2 \left( \frac{Q_1}{4\pi\epsilon_0 \sqrt{1^2}} \right) + \frac{Q_3}{4\pi\epsilon_0} \left[ \frac{Q_1}{\sqrt{1^2}} + \frac{Q_2}{\sqrt{1^2 + 1^2}} \right]$$

$$W = \frac{1}{4\pi\epsilon_0} \left( \frac{Q_1 Q_2}{1} + \frac{Q_1 Q_3}{1} + \frac{Q_2 Q_3}{\sqrt{2}} \right)$$

$$W = 9 \times 10^9 \left[ (-1 \times 4) + (-1 \times 3) + \frac{(4 \times 3)}{\sqrt{2}} \right] \times 10^{-18}$$

$$W = 9 \times 10^9 \left[ -\frac{12}{\sqrt{2}} - 7 \right] = 13.37 \text{ nJ}$$

$$W = \frac{1}{2} \sum_{k=1}^3 Q_k V_k = \frac{1}{2} (Q_1 V_1 + Q_2 V_2 + Q_3 V_3)$$

$$IN = \frac{1}{2} [Q_1(V_{12} + V_{13}) + Q_2(V_{21} + V_{23}) + Q_3(V_{31} + V_{32})]$$

$$W = \frac{Q_1}{2} \left[ \frac{Q_2}{4\pi\epsilon_0(1)} + \frac{Q_3}{4\pi\epsilon_0(1)} \right] + \frac{Q_2}{2} \left[ \frac{Q_1}{4\pi\epsilon_0(1)} + \frac{Q_3}{4\pi\epsilon_0(2)} \right] + \frac{Q_3}{2} \left[ \frac{Q_1}{4\pi\epsilon_0(1)} + \frac{Q_2}{4\pi\epsilon_0(2)} \right]$$

$$W = \frac{1}{4\pi\epsilon_0} \left[ Q_1 Q_2 + Q_1 Q_3 + \frac{Q_2 Q_3}{\sqrt{2}} \right] = \frac{1}{4\pi\epsilon_0} \left[ -4 - 3 + \frac{12}{\sqrt{2}} \right] \times 10^{18}$$

$$W = 9 \times 10^9 \times 10^{18} \left( \frac{12}{\sqrt{2}} - 7 \right) = 13.37 \text{ nJ}$$

Problem (3.1a) Point charges  $Q_1 = 1 \text{ nc}$ ,  $Q_2 = -2 \text{ nc}$ ,  $Q_3 = 3 \text{ nc}$  and  $Q_4 = -4 \text{ nc}$  are positioned one at a time and in that order at  $(0, 0, 0)$ ,  $(1, 0, 0)$ ,  $(0, 0, -1)$  and  $(0, 0, 1)$  respectively. Calculate the energy in the system after each charge is positioned.  
 (Ans: 0,  $-18 \text{ nJ}$ ,  $-29.18 \text{ nJ}$ ,  $-68.27 \text{ nJ}$ )

Line

:  $\overrightarrow{poor}$

## Capacitance:

\* Y V N R \*

EF-18

Capacitor have two conductors carrying equal but opposite charges. The conductors are sometimes referred to as the plates of the capacitor. The plates may be separated by free spaces or a dielectric.

- Consider two conductor capacitor as shown in fig.
- The conductors are maintained at a potential difference  $V$  given by

$$V = V_2 - V_1 = - \int_{\text{2}}^{\text{1}} E \cdot d\ell$$

where  $E$  = Electric field existing between the conductors

Conductor 1 is assumed to carry a positive charge.

Def: Capacitance of a capacitor can be defined as the ratio of the magnitude of the charge on one of the plate to the potential difference between them. i.e.,

$$\begin{aligned} C &= \frac{Q}{V} = \frac{\oint D \cdot dS}{-\int E \cdot d\ell} \\ &= -\frac{\oint \epsilon_0 E \cdot dS}{\int E \cdot d\ell} \end{aligned}$$

$$\therefore C = \epsilon \left| \frac{\oint E \cdot dS}{\int E \cdot d\ell} \right| \quad (8)$$

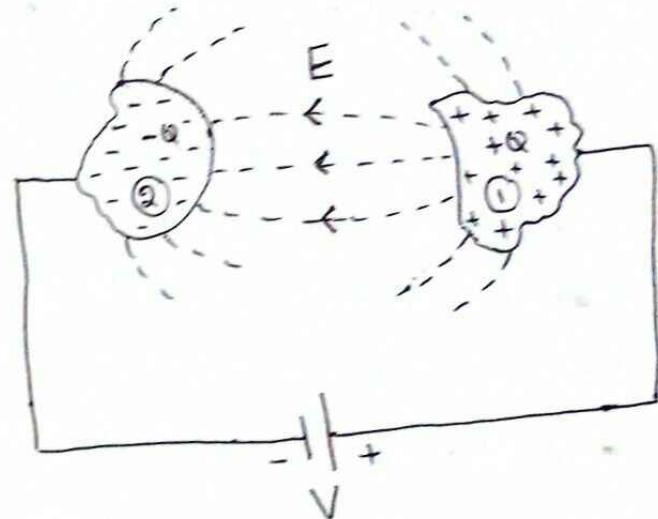


fig. Two conductor capacitor

\* The negative sign before  $V = - \int E \cdot d\ell$  has been dropped because we are interested in the absolute value of  $V$ .

$$C = \frac{\epsilon \oint E \cdot dS}{\int E \cdot d\ell}$$

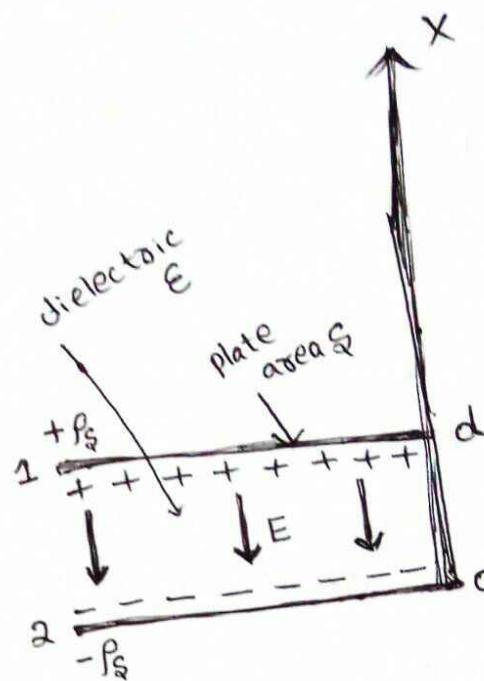
→ Capacitor is denoted with C and is measured in farads (F).

→ There are three types of capacitors.

1. Parallel-plate capacitor
2. Coaxial capacitor
3. Spherical capacitor.

### 1. Parallel-plate capacitor :

Consider a parallel-plate capacitor is shown in fig.  
Suppose that each of the plates has an area  $S$  and they are separated by a distance  $d$ .



→ The space between the plates is filled with a dielectric of permittivity  $\epsilon$ .

→ The plate 1 carries the positive charge and is distributed over its surface with a charge density  $+ρ_s$ .

→ The plate 2 carries the negative charge and is distributed over its surface with a charge density  $-ρ_s$ .

$$\therefore ρ_s = \frac{Q}{S} = \frac{Q}{S}$$

$$\text{The Surface charge } ρ_s = \frac{Q}{S} \quad \longrightarrow (1)$$

→ To find the potential difference, let us obtain  $E$  between the plates.

Assuming plate 1 to be infinite sheet of charge.

$$E_1 = \frac{ρ_s}{2ε} \text{ an} = \frac{ρ_s}{2ε} (-ax) \quad \text{v/m}$$

while for plate 2, we can write

$$E_2 = \frac{-ρ_s}{2ε} \text{ an} = \frac{-ρ_s}{2ε} ax \quad \text{v/m}$$

→ The electric field existing between the two plates having equal and opposite charges is given by

$$\therefore E = E_1 + E_2$$

$$\therefore E = E_1 + E_2$$

$$E = \frac{P_S}{2\epsilon} (-ax) + \frac{-P_S}{2\epsilon} ax$$

$$E = -\frac{P_S}{2\epsilon} ax - \frac{P_S}{2\epsilon} ax$$

$$E = -\frac{2P_S}{2\epsilon} ax$$

$$E = -\frac{P_S}{\epsilon} ax \quad \longrightarrow (2)$$

Substitute eq(1) in (2). Then

$$E = -\frac{Q}{\epsilon S} ax$$

→ The potential difference is given by

$$V = - \int_0^d E \cdot dl$$

$$V = - \int_0^d \left( -\frac{Q}{\epsilon S} ax \right) \cdot dx ax$$

$$V = \frac{Q}{\epsilon S} \int_0^d dx$$

$$V = \frac{Q}{\epsilon S} [x]_0^d$$

$$V = \frac{Qd}{\epsilon S}$$

differential length in  
Cartesian coordinate

System:  $dx ax + dy ay + dz az$

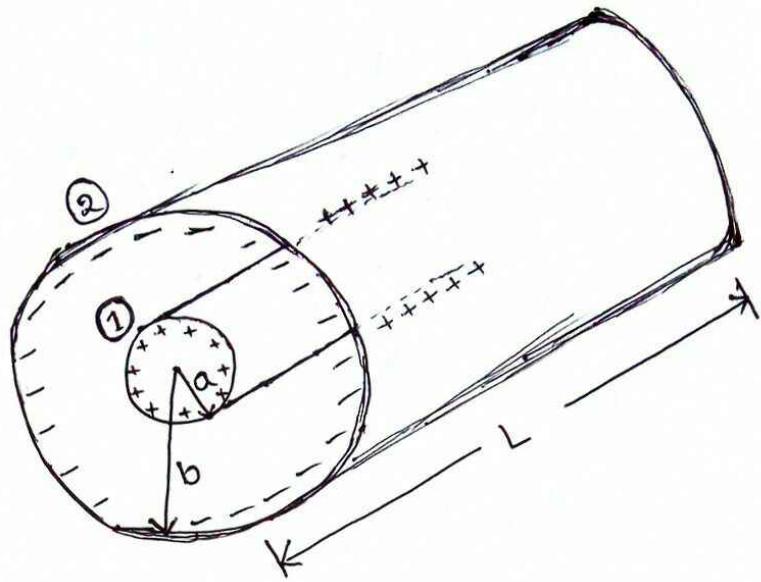
→ The capacitance is the ratio of charge Q to voltage V.

$$C = \frac{Q}{V} = \frac{Q}{Qd/\epsilon S} = \frac{\epsilon S}{d}$$

$\therefore C = \frac{\epsilon S}{d}$

## 2. Coaxial capacitor :

This is essentially a coaxial cable (or) coaxial cylindrical capacitor.  
Consider a length  $L$  of two coaxial conductors of inner radius  $a$  and outer radius  $b$  ( $b > a$ ) as shown in fig.



- let the space between the conductors be filled with a homogeneous dielectric with permittivity  $\epsilon$ .
- The inner conductor 1 carries the positive charge and is distributed along the line with a charge density  $+P_l$
- The outer conductor 2 carries the negative charge and is distributed along the line with a charge density  $-P_l$

$$\text{The line charge } P_l = \frac{Q}{L}$$

$$P_l = \frac{dQ}{dl} = \frac{Q}{L}$$

- Assuming cylindrical coordinate System, Electric field ( $E$ ) will be radial from inner to outer conductor
- \* Use Gauss's law to obtain Electric field b/w inner & outer conductors.

$$\Psi = Q_{\text{enclosed}} = \oint_S D \cdot dS$$

$$Q = \oint_S D_{\text{par}} \cdot P d\phi dz d\phi$$

$$Q = \oint_S D_p P d\phi dz \quad (1)$$

$$Q = P D_p \int_{z=0}^L dz \int_{\phi=0}^{2\pi} d\phi$$

$$Q = P D_p \{z\}_0^L [\phi]_0^{2\pi}$$

In cylindrical system  
 $D = D_p a \hat{r}$   
 $dS = P d\phi dz a \hat{r}$   
 $dl = d\phi a \hat{\theta} + dz a \hat{z}$

$$Q = P D p L 2\pi$$

$$Dp = \frac{Q}{2\pi PL}$$

$$\text{Hence } D = D_p a_p = \frac{Q}{2\pi PL} a_p$$

$$E = \frac{D}{\epsilon_0}$$

$$E = \frac{Q}{2\pi\epsilon_0 PL} a_p$$

→ The Potential difference is given by

$$\begin{aligned} V &= - \int^1 E \cdot dl \\ &= - \int_b^a \frac{Q}{2\pi\epsilon_0 PL} a_p \cdot dP a_p \\ &= - \frac{Q}{2\pi\epsilon_0 L} \int_b^a \frac{1}{P} dP \quad (1) \\ &= - \frac{Q}{2\pi\epsilon_0 L} [\log P]_b^a \\ &= - \frac{Q}{2\pi\epsilon_0 L} [\log a - \log b] \\ &= \frac{Q}{2\pi\epsilon_0 L} [\log b - \log a] \\ &= \frac{Q}{2\pi\epsilon_0 L} \log(b/a) \end{aligned}$$

$P$  → Gaussian cylindrical surface of radius  
 $a < P < b$

$$\int \frac{1}{x} dx = \log x$$

→ The capacitance is the ratio of charge  $Q$  to voltage  $V$

$$C = \frac{Q}{V} = \frac{Q}{\frac{Q}{2\pi\epsilon_0 L} \log(b/a)} = \frac{2\pi\epsilon_0 L}{\log(b/a)}$$

$$\therefore C = \frac{2\pi\epsilon_0 L}{\log(b/a)}$$

\* Note: log is replaced by in in this derivation.

### 3. Spherical Capacitor:

This is the case of two concentric spherical conductors.

Consider the inner sphere of radius 'a' and outer sphere of radius 'b' ( $b > a$ ) separated by a dielectric medium with permittivity  $\epsilon$  as shown in fig.

→ We assume charges  $+Q$  and  $-Q$  on the inner and outer spheres respectively.

→ Use Gauss's law to obtain Electric field b/w inner & outer spheres.

$$\Psi = Q_{\text{enclosed}} = \oint_S D \cdot dS$$

$$Q = \oint D r a \sigma \cdot r^2 \sin \theta d\theta d\phi dr$$

$$Q = D r \sigma^2 \int_{\phi=0}^{2\pi} d\phi \int_{\theta=0}^{\pi} \sin \theta d\theta \quad (1)$$

$$Q = D r \sigma^2 [\phi]_0^{2\pi} [-\cos \theta]_0^{\pi}$$

$$Q = D r \sigma^2 (2\pi) \cdot [ -(-1-1) ]$$

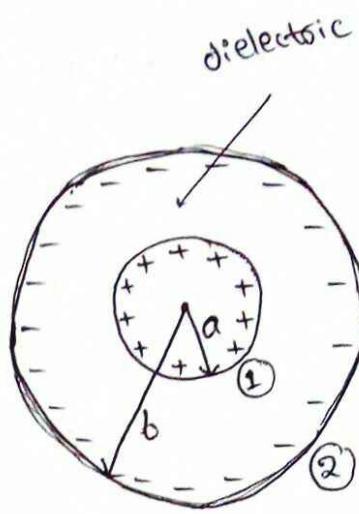
$$Q = D r 4\pi r^2$$

$$D_r = \frac{Q}{4\pi r^2}$$

$$\text{Hence } D = D_r a \sigma = \frac{Q}{4\pi r^2} a \sigma$$

$$E = \frac{D}{\epsilon_0}$$

$$E = \frac{Q}{4\pi \epsilon_0 r^2} a \sigma$$



In spherical system

$$D = D_r a \sigma$$

$$dS = r^2 \sin \theta d\theta d\phi dr$$

$$dl = dr a \sigma + r d\theta a \sigma + r \sin \theta d\phi a \sigma.$$

→ The Potential difference is given by

$$V = - \int E \cdot d\ell$$

$$V = - \int_b^a \frac{Q}{4\pi\epsilon_0 r^2} dr$$

$$V = - \int_b^a \frac{Q}{4\pi\epsilon_0 r^2} dr \quad (1)$$

$$V = - \frac{Q}{4\pi\epsilon_0} \int_b^a \frac{1}{r^2} dr$$

$$V = \frac{-Q}{4\pi\epsilon_0} \left[ -\frac{1}{r} \right]_b^a$$

$$V = \frac{Q}{4\pi\epsilon_0} \left[ \frac{1}{r} \right]_b^a$$

$$V = \frac{Q}{4\pi\epsilon_0} \left[ \frac{1}{a} - \frac{1}{b} \right]$$

$r \rightarrow$  Gaussian spherical  
Surface of radius

$$a < r < b$$

→ The capacitance is the ratio of charge Q to voltage V

$$C = \frac{Q}{V}$$

$$C = \frac{Q}{\frac{Q}{4\pi\epsilon_0} \left[ \frac{1}{a} - \frac{1}{b} \right]}$$

$$\boxed{\therefore C = \frac{4\pi\epsilon_0}{\frac{1}{a} - \frac{1}{b}}}$$

Problem (6.13) A cylindrical capacitor has radii  $a=1\text{cm}$  &  $b=2.5\text{cm}$ . If the space b/w the plates is filled with an inhomogeneous dielectric with  $\epsilon_r = (v+p)/p$ , where  $p$  in centimeters. find the capacitance per meter of the capacitor.

Sol Electric field due to line charge  $E = \frac{\alpha L}{2\pi \epsilon_0 p} = \frac{\alpha}{2\pi \epsilon_0 p L} \quad (\because R = \frac{\alpha}{L})$

$$\therefore V = - \int_b^a \frac{\alpha}{2\pi \epsilon_0 p L} dp = \frac{-\alpha}{2\pi \epsilon_0 L} \int_b^a \frac{dp}{p(10+p)} = \frac{-\alpha}{2\pi \epsilon_0 L} \int_b^a \frac{dp}{10+p} \quad [\because v = - \int E \cdot dL] \quad [L=1]$$

$$V = \frac{-\alpha}{2\pi \epsilon_0 L} \left[ \ln(10+p) \right]_b^a = \frac{\alpha}{2\pi \epsilon_0 L} \ln \left( \frac{10+b}{10+a} \right) = \frac{\alpha}{2\pi \epsilon_0} \ln \left( \frac{12.5}{11} \right) = \frac{\alpha}{2\pi \epsilon_0} \ln(0.127)$$

$$C = \frac{\alpha}{V} = \frac{2\pi \epsilon_0}{\ln(0.127)} = \frac{2\pi \cdot 10^{-9}}{36\pi \ln(0.127)} = 434.6 \text{ PF} \quad \boxed{\therefore C = 434.6 \text{ PF}}$$

Problem A spherical capacitor with  $a=1.5\text{cm}$  &  $b=4\text{cm}$  has an inhomogeneous dielectric of  $\epsilon = \frac{10\epsilon_0}{r}$ . calculate the capacitance of capacitor Ans: 1.13 nF

Problem Given the electric flux density  $D = 0.3\alpha r \text{ a.s nc/m}^2$  in free space (i) Find  $E$  at  $r(=2, \theta=25^\circ, \phi=90^\circ)$ . (ii) Find the total charge within the sphere of radius  $r=3$  (iii) Find the total electric flux leaving the sphere of radius  $r=4$ .

Sol (i)  $D = \epsilon_0 E \Rightarrow E = \frac{D}{\epsilon_0} = \frac{(0.3)\alpha r}{\epsilon_0} \text{ a.s nc/m}^2 = \frac{0.3(2) \alpha r \times 10^{-9}}{10^{-9}/36\pi} = 0.6 \times 36\pi = 67.85 \text{ V/m}$

(ii)  $Q = \oint D \cdot dS = \oint (0.3)\alpha r \cdot r^2 \sin\theta d\theta d\phi dr = \int_{\phi=0}^{2\pi} d\phi \cdot \int_{\theta=0}^{\pi} \sin\theta d\theta \cdot (0.3) r^3 \quad (1)$

$$Q = \left[ \phi \right]_0^{2\pi} \left[ -\cos\theta \right]_0^{\pi} (0.3)(3)^3 = 2\pi \cdot 2 \cdot (0.3) \cdot (27) = 101.78 \text{ nc}$$

(iii)  $Q = \Psi \quad (r=4) :$

$$Q = D(4\pi r^2) = (0.3)r \cdot \alpha r(4\pi r^2) = (0.3)4\pi r^3$$

$$Q = (0.3)4\pi(4)^3 = 241.27 \text{ nc}$$

$$\begin{aligned} \frac{Q}{4\pi r^2} &= D \\ D &= D \alpha r \\ D &= D \cdot r \end{aligned}$$

Problem Find the total charge inside each of the volume indicated

(i)  $P_V = 10 z^2 e^{0.1x} (\sin\pi y), -1 \leq x \leq 2, 0 \leq y \leq 1, 3 \leq z \leq 3.6$

(ii)  $P_V = 4xyz^2, 0 \leq x \leq 2, 0 \leq y \leq \pi/2, 0 \leq z \leq 3$

(iii)  $P_V = \frac{3\pi \cos^2 \theta \cos^2 \phi}{2r^2(r^2+1)}$